

Impact of complex phases on MSSM Higgs searches

Interference and other \mathcal{CP} -violating effects.

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DESY

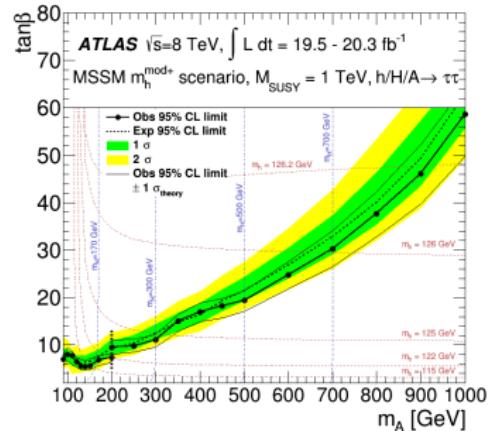
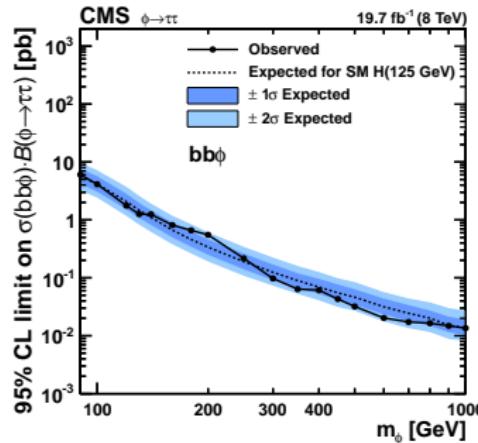
in collaboration with Sven Heinemeyer,
Oscar Stål and Georg Weiglein

Terascale Workshop, Hamburg
December 2, 2014
Higgs session

Interpretation of MSSM Higgs searches

Experimental searches for $\Phi = h, H, A \rightarrow \tau^+\tau^-, \mu^+\mu^-, b\bar{b}$ → A. Gilbert's, L. Hauswald's talks

production $\{gg \rightarrow \Phi, b\bar{b}\Phi\} \times$ decay $\Phi \rightarrow \{\tau^+\tau^-, \mu^+\mu^-, b\bar{b}\}$



Limitation of factorisation in standard NWA

interference terms neglected, relevant especially with complex phases



Outline

- 1 Motivation**
- 2 Complex phases in the MSSM Higgs sector**
- 3 Interference effects**
- 4 Phenomenological application**
- 5 Conclusion**



Complex phases: motivation and constraints

Motivation

- ▶ baryon asymmetry of the universe requires more \mathcal{CP} -violation than in CKM matrix
- ▶ **MSSM Higgs sector is \mathcal{CP} -conserving at lowest order**
- ▶ parameters from **other sectors** can in principle be **complex**: 12
 - trilinear couplings A_f
 - higgsino mass parameter μ
 - gaugino mass parameters M_1, M_2 (rotate ϕ_{M_2} away), M_3



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Constraints from EDMs (Tl, Hg, n, D)

e.g. [Barger, Falk, Han, Jiang, Li, Plehn '01], [Ellis, Lee, Pilaftsis '09], [Li, Profumo, Ramsey-Musolf '10]

- ▶ $\phi_{A_{f1,2}}$ more strongly constrained than $\phi_{A_{t,b}}$
- ▶ ϕ_{M_1} can be sizeable - large
- ▶ ϕ_{M_3} strongly constrained only if $\tilde{f}_{1,2}$ light
- ▶ ϕ_μ tight limits

Most relevant phases in Higgs sector: $\phi_{A_{t,b}}, \phi_{M_3}$, enhanced $\mu \cdot A_t$



Real case: \mathcal{CP} conserved

- ▶ only \mathcal{CP} -even states mix:
 $h - H$
- ▶ M_A or M_{H^\pm} as input mass

Complex case: \mathcal{CP} violated

- ▶ all neutral states mix:
 $h, H, A \rightarrow h_1, h_2, h_3$
- ▶ M_{H^\pm} as input mass

\mathcal{CP} -violating phases can cause interesting phenomenology

Criteria for a significant interference effect

1.) Degeneracy

Nearby resonances

- ▶ masses M_i, M_j
- ▶ widths Γ_i, Γ_j

overlap if $\Delta M \leq \Gamma_i, \Gamma_j$

2.) Mixing

- ▶ Matrix elements $\mathcal{M}_i, \mathcal{M}_j$
- ▶ if i, j do not mix:
$$\sigma_{\text{Int}} \propto 2\text{Re}[\mathcal{M}_i \mathcal{M}_j^*] = 0$$



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Interference in MSSM Higgs sector?

- ▶ real parameters: only h, H mix
 - but $M_h \simeq M_H$ limited to narrow parameter range
- ▶ complex parameters: all neutral Higgs bosons mix $\rightarrow h_1, h_2, h_3$
 - $M_{h_3} - M_{h_2} \leq \Gamma_{h_2}, \Gamma_{h_3}$ in decoupling region

Analyse interference effects between neutral Higgs bosons!



\mathcal{CP} -violating Higgs interference

In presence of non-zero phase: change of cross section

- ▶ dominant effect: $h_2 - h_3$ interference $\Rightarrow \sigma_{H+A} \not\approx 2\sigma_H$ or $\sum \sigma_\Phi \text{BR}_\Phi$
- ▶ also affected: couplings, masses, widths, mixings in/outside σ_{int}

Our approach

- ▶ full propagator mixing $\Delta_{ij}(p^2)$: 3×3 or 2×2 (no NWA)
 - $\phi \equiv \phi_{A_t} \neq 0$ or $\phi = 0 \rightarrow \boxed{\delta := \frac{\sigma(\phi) - \sigma(0)}{\sigma(0)}}$
- ▶ measures relative effect of complex phase on cross section σ



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Process: $b\bar{b} \rightarrow \tau^+\tau^-$ via $\Phi = h_1, h_2, h_3$

$$\left| \overline{b} \begin{array}{c} b \\ \diagup \\ h_1 \\ \diagdown \\ \tau^- \end{array} + \overline{b} \begin{array}{c} b \\ \diagup \\ h_2 \\ \diagdown \\ \tau^- \end{array} + \overline{b} \begin{array}{c} b \\ \diagup \\ h_3 \\ \diagdown \\ \tau^- \end{array} \right|^2$$

FeynArts,
FormCalc,
FeynHiggs



Benchmark scenario: $M_h^{\text{mod+}}$

$$M_{\text{SUSY}} = 1000 \text{ GeV}$$

$$M_2 = 200 \text{ GeV}$$

$$X_t^{\text{OS}} = 1.5 M_{\text{SUSY}}$$

$$A_t = A_b = A_\tau$$

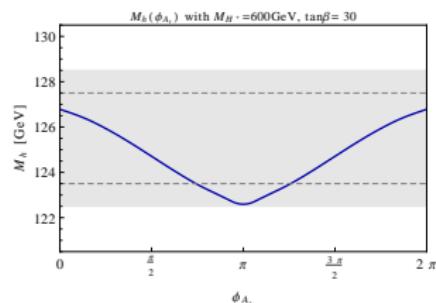
$$M_3 = 1500 \text{ GeV}$$

$$M_{\tilde{f}_3} = M_{\text{SUSY}}$$

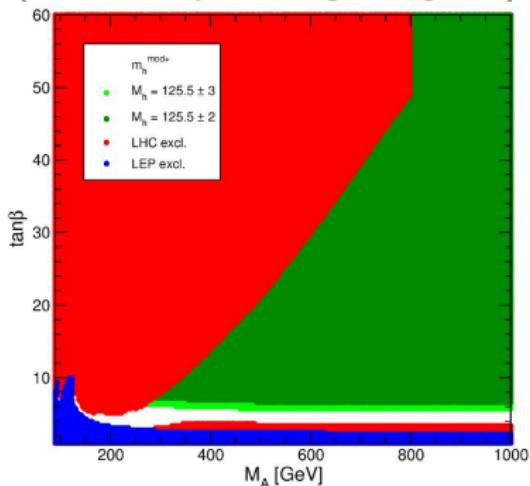
$$M_{\tilde{q}_{1,2}} = 1500 \text{ GeV}$$

$$M_{\tilde{l}_{1,2}} = 500 \text{ GeV}$$

$$\mu = \pm 200, \pm 500, \pm 1000 \text{ GeV}$$



[Carena, Heinemeyer, Stål, Wagner, Weiglein '13]



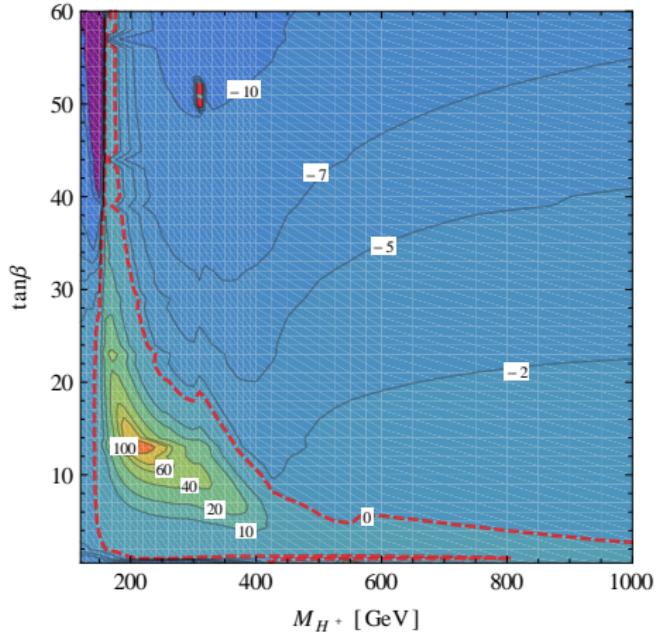
- ▶ most parameter space M_h -allowed
- ▶ ϕ_{A_t} changes M_h
- ▶ but M_h mostly stays in allowed region

Impact of ϕ_{A_t} in $M_h^{\text{mod}+}$ scenario

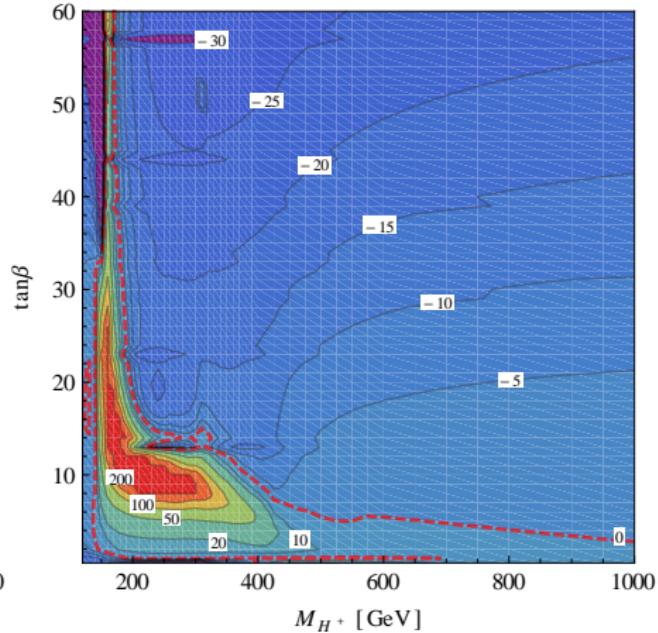
preliminary

$$\delta = \frac{\sigma_\phi - \sigma_0}{\sigma_0} \text{ in full plane}$$

$M_h^{\text{mod}+}$, $\mu = 200$ GeV: $\delta(\phi_{A_t} = \pi/4)$ [%]



$M_h^{\text{mod}+}$, $\mu = 200$ GeV: $\delta(\phi_{A_t} = \pi/2)$ [%]



Constructive and destructive effects of δ

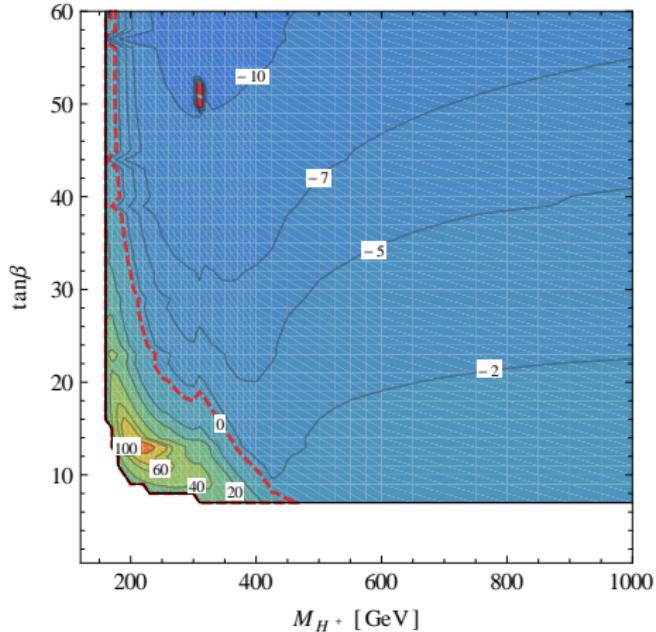


Impact of ϕ_{A_t} in $M_h^{\text{mod}+}$ scenario

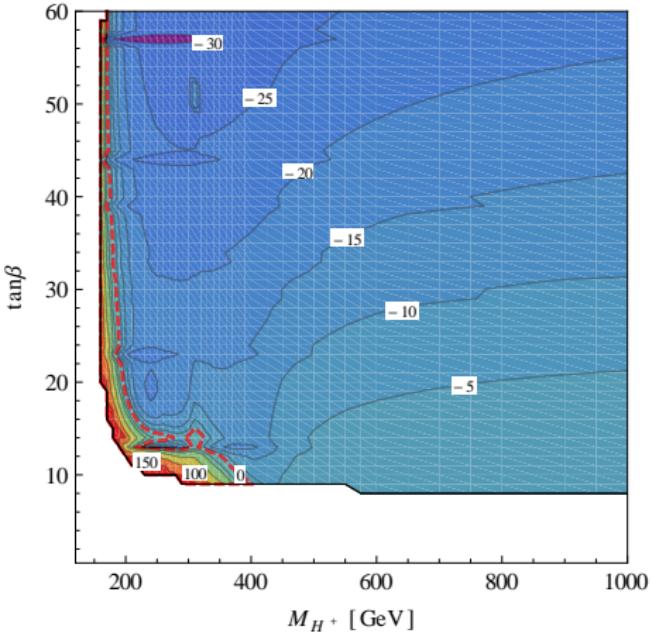
preliminary

$$\delta = \frac{\sigma_\phi - \sigma_0}{\sigma_0} \text{ where } M_h = 125.5 \pm 3 \text{ GeV}$$

$M_h^{\text{mod}+}, \mu = 200 \text{ GeV}: \delta(\phi_{A_t} = \pi/4)[\%]$



$M_h^{\text{mod}+}, \mu = 200 \text{ GeV}: \delta(\phi_{A_t} = \pi/2)[\%]$



Mostly destructive effects of δ in M_h -allowed region



back-of-the-envelope calculation:

- ▶ for fixed M_A
- ▶ $\sigma(0, \tan \beta_0) \leftrightarrow$ exclusion limit
- ▶ $\sigma(\phi, \tan \beta_0) = (1 + \delta)\sigma(0, \tan \beta_0)$



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- ▶ adjust $\tan \beta_0 \mapsto \tan \beta_\phi$ s.t.
 $\sigma(\phi, \tan \beta_\phi) = \sigma(0, \tan \beta_0)$



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 $\sigma(\phi, \tan \beta_\phi) = \sigma(0, \tan \beta_0)$
- ▶ $\sigma(\tan \beta) \propto \tan \beta^2$
- ▶
$$\tan \beta_\phi \simeq \frac{\tan \beta_0}{\sqrt{(1 + \delta)}}$$



back-of-the-envelope calculation:

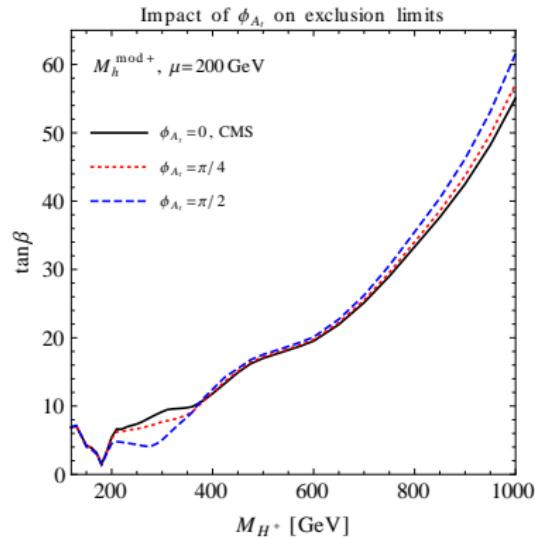
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- ▶ $\sigma(\tan \beta) \propto \tan \beta^2$
- ▶
$$\tan \beta_\phi \simeq \frac{\tan \beta_0}{\sqrt{(1 + \delta)}}$$

More precise $\tan \beta$ -limit

- ▶ HiggsBounds

Comparison with CMS

24.6/fb, [1408.3316], $\tau^+ \tau^-$



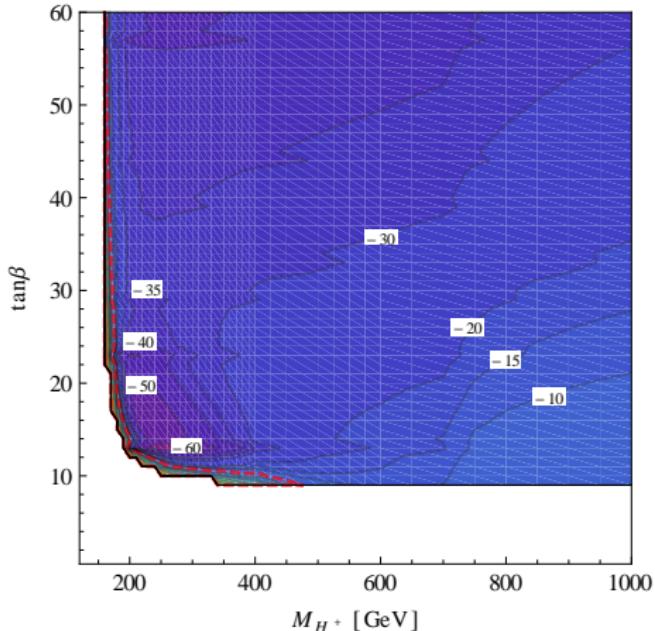
slightly weakened limit at
high masses for $\phi_{A_t} \neq 0$

Impact of μ in $M_h^{\text{mod+}}$ scenario

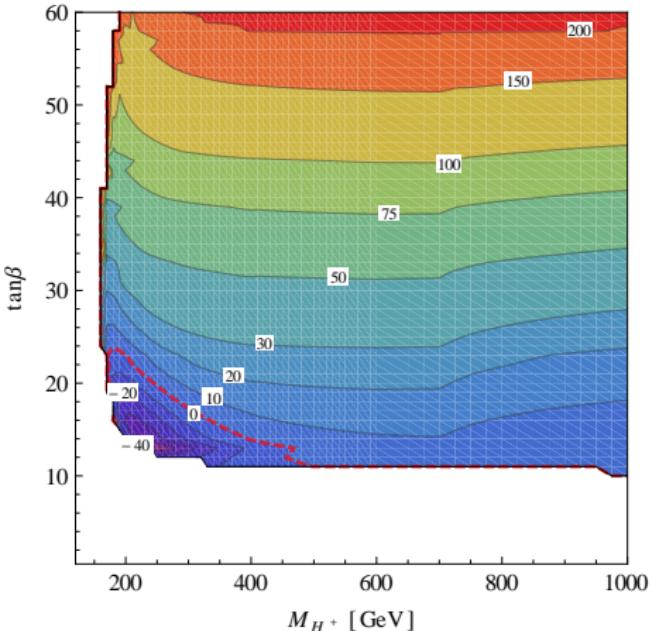
preliminary

$$\mu = \pm 500 \text{ GeV}: \delta = \frac{\sigma_\phi - \sigma_0}{\sigma_0} \text{ where } M_h = 125.5 \pm 3 \text{ GeV}$$

$M_h^{\text{mod+}} \mu = +500 \text{ GeV}: \delta [\%]$ effect of $\phi_{A_t} = \pi/2$



$M_h^{\text{mod+}} \mu = -500 \text{ GeV}: \delta [\%]$ effect of $\phi_{A_t} = \pi/2$

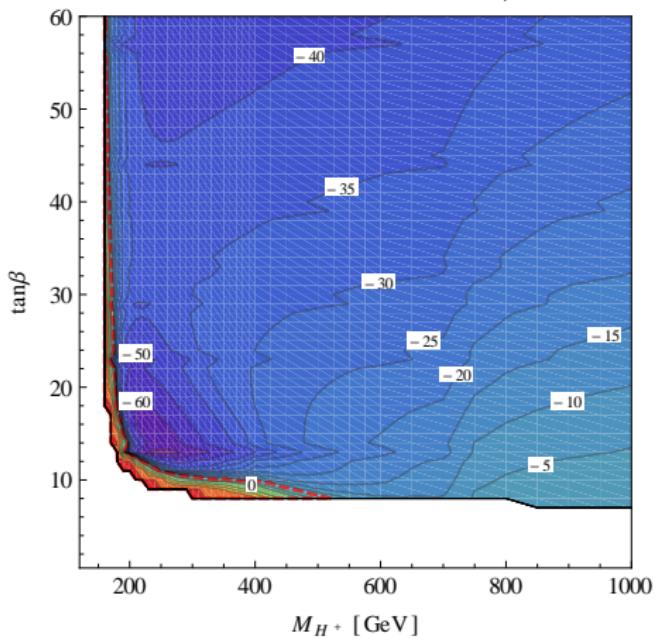


$|\mu|, \text{sgn}(\mu)$ have crucial influence on δ

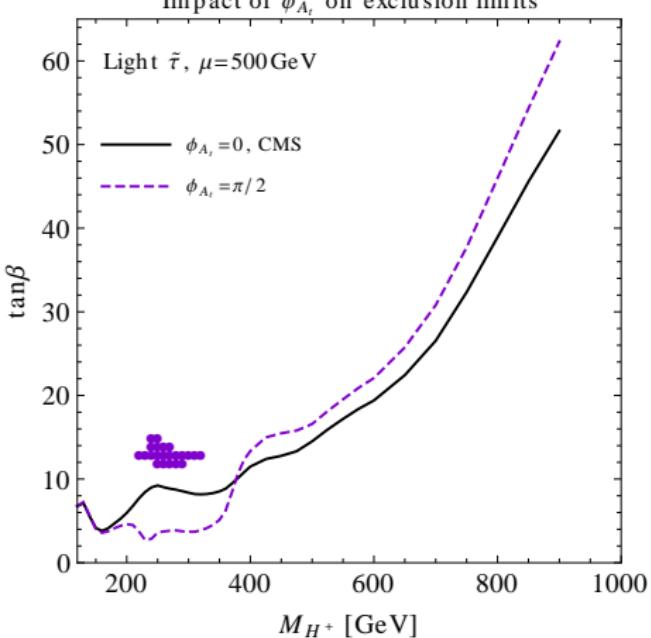
$\mu = 500 \text{ GeV}$ in light- $\tilde{\tau}$ scenario

preliminary

Light $\tilde{\tau}$ scenario : Relative effect of $\phi_{A_t} = \pi/2$



Impact of ϕ_{A_t} on exclusion limits



Unexcluded hole of parameter points due to
strongly destructive interference effect

Conclusion: relevant phase effects

Outlook

- ▶ Calculate the impact of ϕ_{At} and interference also in gluon fusion
- ▶ Study impact on experimental limits in more detail: HiggsBounds
- ▶ Separate interference term from other non-zero phase effects



Conclusion: relevant phase effects

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- ▶ Calculate the impact of ϕ_{At} and interference also in gluon fusion
- ▶ Study impact on experimental limits in more detail: HiggsBounds
- ▶ Separate **interference term** from other **non-zero phase** effects

Summary

- ▶ Complex phases allow for \mathcal{CP} -violating mixing and interference effects in the MSSM Higgs sector
 - \mathcal{CP} -violating case: analysis beyond *incoherent* sum $\sigma_A \text{BR}_A + \sigma_H \text{BR}_H$
- ▶ Extend analyses of benchmark scenarios
 - include different values of $\mu = \pm 200, \pm 500, \pm 1000 \text{ GeV}$
- ▶ Non-zero phases can have **significant impact on exclusion limits**



Conclusion: relevant phase effects

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- ▶ Calculate the impact of ϕ_{At} and interference also in gluon fusion
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Thank you!



The MSSM Higgs sector

2 Higgs doublets needed for holomorphic superpotential

$$\mathcal{H}_1 = \begin{pmatrix} H_{11} \\ H_{12} \end{pmatrix} = \begin{pmatrix} v_1 + \frac{1}{\sqrt{2}}(\phi_1^0 - i\chi_1^0) \\ -\phi_1^- \end{pmatrix}$$
$$\mathcal{H}_2 = \begin{pmatrix} H_{21} \\ H_{22} \end{pmatrix} = \begin{pmatrix} \phi_2^+ \\ v_2 + \frac{1}{\sqrt{2}}(\phi_2^0 + i\chi_2^0) \end{pmatrix} e^{i\xi}$$

Higgs potential with SUSY and soft SUSY terms

$$V_H = m_1^2 \mathcal{H}_1^\dagger \mathcal{H}_1 + m_2^2 \mathcal{H}_2^\dagger \mathcal{H}_2 - m_{12}^2 \epsilon^{ij} (\mathcal{H}_{1i} \mathcal{H}_{2j} + \text{h.c.})$$
$$+ \frac{1}{8} (g_1^2 + g_2^2) (\mathcal{H}_1^\dagger \mathcal{H}_1 - \mathcal{H}_2^\dagger \mathcal{H}_2)^2 + \frac{1}{2} g_2^2 |\mathcal{H}_1^\dagger \mathcal{H}_2|^2$$

$$m_i^2 = \tilde{m}_i^2 + |\mu|^2, \quad i = 1, 2$$

$$m_{12}^2 = b\mu = |m_{12}^2| e^{i\xi}$$

Complex phases ξ, ζ can be rotated away



Relations at lowest order

Higgs sector \mathcal{CP} -conserving at lowest order

Physical states

- \mathcal{CP} -even: $\phi_1^0, \phi_2^0 \rightarrow h, H$
- \mathcal{CP} -odd: $\chi_1^0, \chi_2^0 \rightarrow A, G$
- charged: $\phi_1^\pm, \phi_2^\pm \rightarrow H^\pm, G^\pm$

Input parameters

$$\tan \beta = \frac{v_2}{v_1}, \quad m_A^2 = \frac{2|m_{12}^2|}{\sin(2\beta)}$$

Masses determined by $m_A, \tan \beta$

$$m_{h/H}^2 = \frac{1}{2} \left(m_A^2 + M_Z^2 \mp \sqrt{(m_A + M_Z^2)^2 - 4m_A^2 M_Z^2 \cos^2(2\beta)} \right)$$

$$\Rightarrow m_h^2 \leq M_Z^2$$

$$m_{H^\pm}^2 = m_A^2 + M_W^2$$

But higher-order corrections important



Higher-order effects in the Higgs sector

Self-energy diagrams

- significant impact on masses

- mixing self-energies $\hat{\Sigma}_{ij}(p^2) \Rightarrow$ mass matrix $\mathbf{M}_{ij} = m_i^2 \delta_{ij} - \hat{\Sigma}_{ij}(p^2)$
- higher-order masses and widths from complex poles of the propagators

$$\mathcal{M}_{h_a}^2 = M_{h_a}^2 - i M_{h_a} \Gamma_{h_a}$$

- particles and parameters from other sectors contribute, e.g.

- sfermions: trilinear couplings A_f and mixing parameters
 $X_{f_{u,d}} = A_{f_{u,d}} - \mu^* \{\cot\beta, \tan\beta\}$ enter at 1-loop order
- gluino: M_3 contribute only at 2-loop order,
but at 1-loop in correction Δ_b to bottom Yukawa coupling

Phenomenology

- loop corrected Higgs mass at 126 GeV possible in MSSM
- important loop contributions to Higgs production and decay

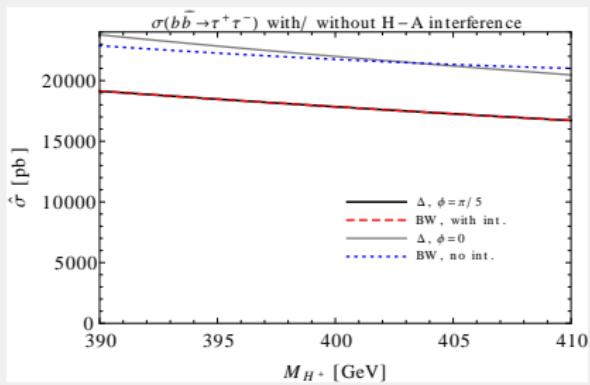
More parameters than $m_A, \tan\beta$ necessary to describe Higgs sector



Complex parameter in $b\bar{b} \rightarrow \tau^+\tau^-$ via $\Phi = h, H, A$

\mathcal{CP} -violating scenario

$\phi_{At} = \pi/5$, $|A_t| = 1200$,
 $\tan \beta = 50$, $M_{\text{SUSY}} = 500 \text{ GeV}$,
 $M_{\tilde{f}} = M_{\text{SUSY}}$, $M_{\tilde{u}_3} = 0.8M_{\text{SUSY}}$
 $\mu = M_2 = 200 \text{ GeV}$



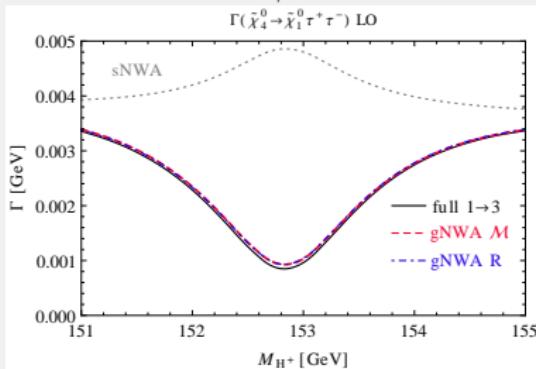
Comparison

- ▶ full cross sections (Δ_{ij}, BW) agree very well
- ▶ phase and interference effect similar
 - significant effects
 - dominant: interference term
- ▶ in the following:
 - full mixing propagators
 - impact of ϕ_{At} on all terms

Example in \mathcal{CP} -conserving Higgs sector

Destructive h-H interference

example process: $\tilde{\chi}_4^0 \rightarrow \tilde{\chi}_1^0 \tau^+ \tau^-$
 \mathbb{R} scenario, $\tan \beta = 50$



sNWA: standard NWA **insufficient here!**

full: 3-body decay BW

gNWA: generalised NWA

Include interference term

Mixing propagators

- full p^2 -dependence
- $\hat{\Sigma}_{ij}$ from FeynHiggs

Breit-Wigner propagators

- approximate p^2 -dependence
- \hat{Z} -factors from FeynHiggs

generalised NWA

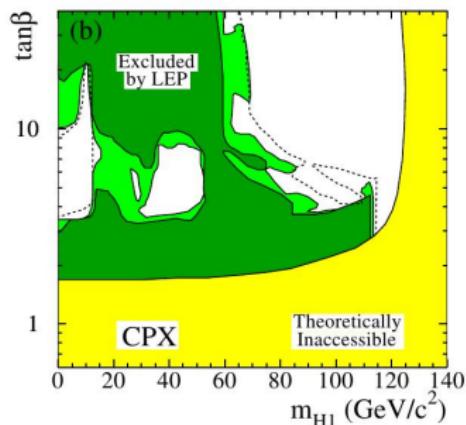
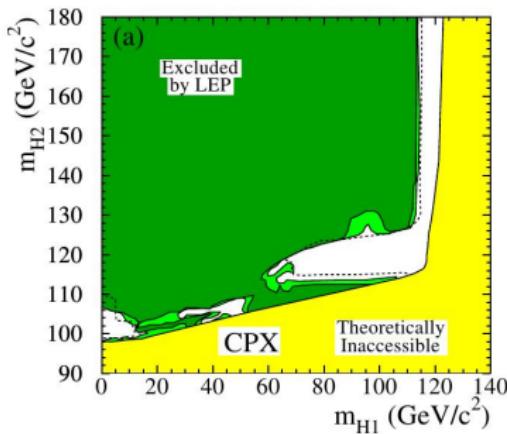
[EF, S. Thewes, G. Weiglein '14]

- on-shell matrix elements
- enables factorisation into production \times decay

Another complex example: CPX hole at LEP

LEP search for neutral MSSM Higgs bosons

interpreted in **CPX scenario** with large amount of \mathcal{CP} -violation



Unexcluded hole in parameter space

due to suppressed coupling of $h_1 ZZ$ in CPX scenario