

Introduction to the Minimal Supersymmetric Standard Model

Ulrich Langenfeld

DESY, Zeuthen

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1 Outline

- Open questions of the Standard Model
- The Particle spectrum of the Minimal Supersymmetric Standard Model
- Examples of SUSY processes
- Squark pair production at the LHC

2 Introduction

The Standard Model (SM) has been tested to high precision.

However, many questions are still open:

- solution to hierarchy problem
($M_W = \mathcal{O}(100 \text{ GeV}) \leftrightarrow M_{\text{Planck}} = \mathcal{O}(10^{19} \text{ GeV})$)
- window to gravity
- dark matter candidate
- CP-phases for baryon asymmetry, baryogenesis
- 19 free parameters in the SM

One solution to some of these questions is Supersymmetry (SUSY).

- symmetry between bosons and fermions
- minimal SUSY extension of SM \rightarrow MSSM

3 The Particle content of the MSSM

Matter fields:

| Fermionic SM-field $s = \frac{1}{2}$ | SUSY partner fields $s = 0$ |
|---|---|
| Electron e_L , Neutrino ν | selectron \tilde{e}_L , sneutrino $\tilde{\nu}$ |
| Electron e_R | selectron \tilde{e}_R |
| up-quark $u_{L,a}$, down-quark $d_{L,a}$ | sup $\tilde{u}_{L,a}$, sdown $\tilde{d}_{L,a}$ |
| up-quark $u_{R,a}$ | sup $\tilde{u}_{R,a}$ |
| down-quark $d_{R,a}$ | sdown $\tilde{d}_{R,a}$ |

$a : SU(3)$ color index

The Particle content of the MSSM

(cont'd)

Gauge fields:

| Bosonic SM-fields $s = 1$ | SUSY partner fields $s = \frac{1}{2}$ |
|---------------------------|---------------------------------------|
| B field | Bino \tilde{B} |
| weak bosons W^a | Winos $\tilde{W}^a, a = 1 \dots 3$ |
| Gluons g^a | Gluinos $\tilde{g}^a, a = 1 \dots 8$ |

The Particle content of the MSSM

(cont'd)

Higgs fields (complex doublet fields):

| Higgs-fields $s = 0$ | SUSY partner fields $s = \frac{1}{2}$ |
|----------------------|--|
| H_d | $\widetilde{H}_1 = (\widetilde{H}_d^0, \widetilde{H}_d^-)$ |
| H_u | $\widetilde{H}_2 = (\widetilde{H}_u^+, \widetilde{H}_u^0)$ |

- Fields with equal quantum numbers mix
- Physical states obtained by diagonalisation of the mixing matrices

Remarks:

- # of fermionic degrees of freedom = # of bosonic d. o. f.
- particles and their super partners have equal masses
- if lightest supersymmetric particle (lsp) is stable $\rightarrow R$ -parity \rightarrow dark matter candidate, if uncharged
- in our world:
no SUSY particles observed so far \Rightarrow SUSY must be broken
- explicit breaking: Introduce mass terms for the the bino, wino, gluino, and sfermions with mass parameters M_1 , M_2 , M_3 and $M_{\tilde{f}}$, respectively
- in SUSY GUTs: $M_1 = M_2 = M_3 = M_{1/2} @ M_{\text{GUT}}$
 $M_{\tilde{f}} = M_0$ for all sfermions $@ M_{\text{GUT}}$

4 Charginos $\tilde{\chi}_{1/2}^{\pm}$

Charginos are the superpartners of the charged gauge \tilde{W}^{\pm} and Higgs bosons $\tilde{H}_{u,d}^0$. Their mixing is described by the matrix

$$X = \begin{pmatrix} M_2 & \sqrt{2}m_W \sin \beta \\ \sqrt{2}m_W \cos \beta & \mu \end{pmatrix}$$

MSSM parameters:

M_2, μ : Wino, higgsino mass parameter, can be complex (CP - phases)

$\tan \beta$: ratio of the Higgs vevs

eigenvalues(X^+X) = chargino masses² $m_{\tilde{\chi}_{1/2}^{\pm}}^2$

Charginos have electromagnetic interactions (good detectability)!

LEP search: $m_{\tilde{\chi}_1^+} > 104 \text{ GeV} \Rightarrow |\mu|, M_2 \gtrsim 100 \text{ GeV}$

5 Neutralinos

Neutralinos are the SUSY partners of the neutral gauge (\tilde{B}, \tilde{W}^3) and CP -even Higgs bosons $(\tilde{H}_u^0, \tilde{H}_d^0)$. These states mix, and the mass eigenstates are the eigenvectors of the diagonalisation matrix.

$$M = \begin{pmatrix} \color{red}{M_1} & 0 & -m_Z \sin(\theta_W) \cos(\color{red}{\beta}) & m_Z \sin(\theta_W) \sin(\color{red}{\beta}) \\ 0 & \color{red}{M_2} & m_Z \cos(\theta_W) \cos(\color{red}{\beta}) & -m_Z \cos(\theta_W) \sin(\color{red}{\beta}) \\ -m_Z \sin(\theta_W) \cos(\color{red}{\beta}) & m_Z \cos(\theta_W) \cos(\color{red}{\beta}) & 0 & -\color{red}{\mu} \\ m_Z \sin(\theta_W) \sin(\color{red}{\beta}) & -m_Z \cos(\theta_W) \sin(\color{red}{\beta}) & -\color{red}{\mu} & 0 \end{pmatrix}$$

MSSM parameters: $\color{red}{M_1}, \color{red}{M_2}, \color{red}{\mu}, \tan \beta$

|eigenvalues| of M = neutralino masses $m_{\tilde{\chi}_i^0}, i = 1, \dots, 4$

- We have some information about the neutralino mixing matrix from the chargino sector ($M_2, |\mu| \geq 100 \text{ GeV}$)
- but one new parameter: M_1 (Bino mass parameter)
- in SUSY GUTs: $M_1 = f(M_2)$,
f. e. mSugra: $M_1 = \frac{5}{3} \tan^2(\theta_w) M_2 \approx 0.5 M_2 @ M_Z$
 $\Rightarrow m_{\tilde{\chi}_1^0} \gtrsim 50 \text{ GeV}$
- no lower bound, if one assumes no relation between M_1 and M_2 , even a massless neutralino is possible [my PhD thesis]
- masses in the range $1 \text{ eV} \dots 6 \text{ GeV}$ disfavoured by cosmology
- $\tilde{\chi}_1^0$ is dark matter candidate: lsp, uncharged, uncoloured

6 Squarks and Sleptons

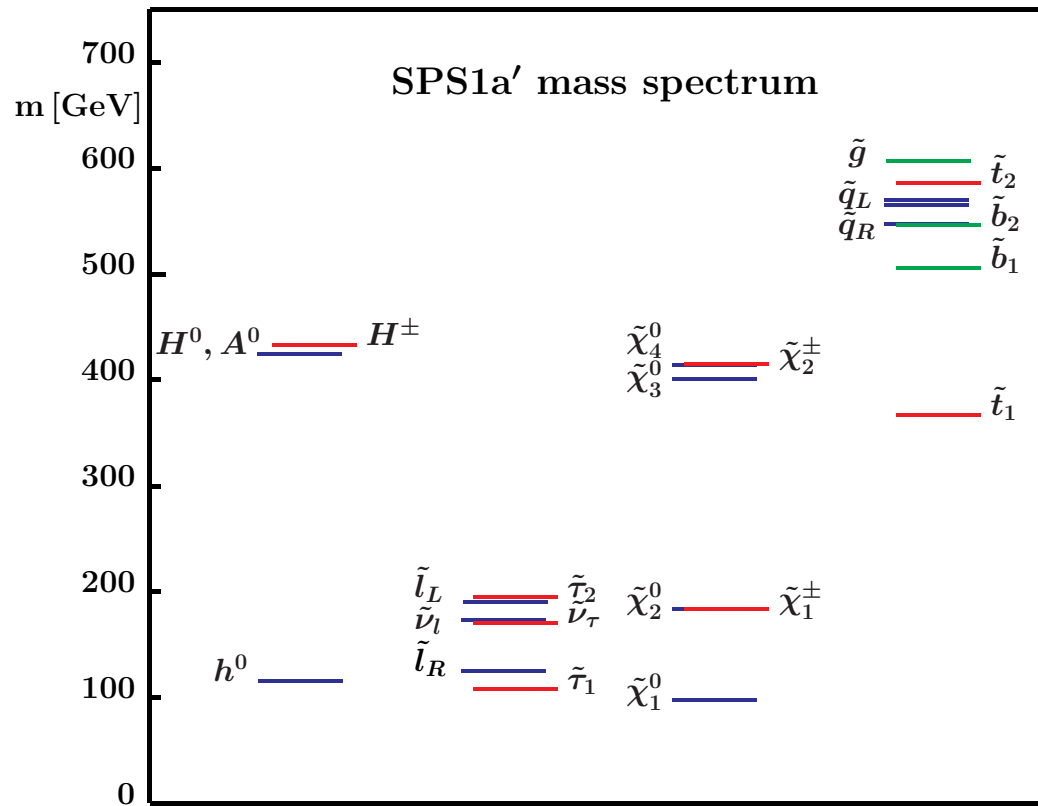
$$M_{\tilde{u}}^2 = \begin{pmatrix} M_{\tilde{q}_L}^2 + m_Z^2 \cos(2\beta)(T_3 - e_u \sin^2 \theta_w) + m_u^2 & m_u(A_u - \mu \cot \beta) \\ m_u(A_u - \mu \cot \beta) & M_{\tilde{u}_R}^2 - e_u m_Z^2 \cos(2\beta) \sin^2 \theta_w + m_u^2 \end{pmatrix}$$

$$M_{\tilde{d}}^2 = \begin{pmatrix} M_{\tilde{q}_L}^2 + m_Z^2 \cos(2\beta)(T_3 - e_d \sin^2 \theta_w) + m_d^2 & m_d(A_d - \mu \tan \beta) \\ m_d(A_d - \mu \tan \beta) & M_{\tilde{d}_R}^2 - e_d m_Z^2 \cos(2\beta) \sin^2 \theta_w + m_d^2 \end{pmatrix}$$

(the matrices for the sleptons look similar)

- $M_{\tilde{u}, \tilde{d}}^2$ is diagonal, if $m_{u,d} = 0 \Rightarrow \tilde{u}_{R,L}, \tilde{d}_{R,L}$ are also mass eigenstates
- in mSugra: Left and right handed Squarks of the 1st and 2nd generation have equal masses \Rightarrow difficult to distinguish at colliders
- Bounds for 1st and 2nd generation: $m_{\tilde{q}} > 379 \text{ GeV} @ 95\% \text{ CL}$ (and $m_{\tilde{g}} > 308 \text{ GeV}$) v. M. Abazov *et al.* [D0 Collaboration], Phys. Lett. B **660** (2008) 449 [arXiv:0712.3805 [hep-ex]].

7 A Typical Particle spectrum



mSugra:

$$M_{1/2} = 250 \text{ GeV}$$

$$M_0 = 70 \text{ GeV}$$

$$A_0 = -300 \text{ GeV}$$

$$\tan \beta = 10$$

$$\text{sgn}(\mu) = +1$$

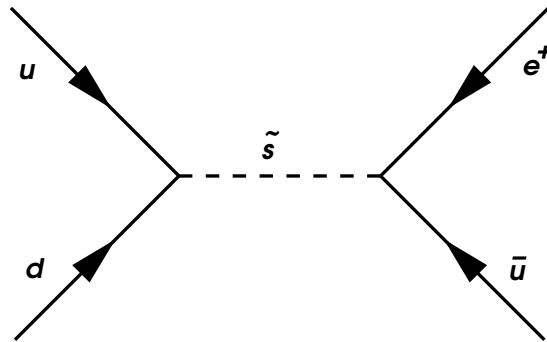
(Figure taken from J. A. Aguilar-Saavedra *et al.*, Eur. Phys. J. C **46** (2006) 43 [arXiv:hep-ph/0511344].)

This scenario fulfills all bounds from HEP and cosmology

8 R -parity

8.1 Motivation and Definition

In a general MSSM, proton decay is possible:



Protect the proton by a new symmetry: R -parity $R_p = (-1)^{3B+L+2S}$

Examples: $R_p(u) = (-1)^{3 \cdot 1/3 + 0 + 2 \cdot 1/2} = 1$ (as for all SM particles),

$R_p(\tilde{e}) = (-1)^{3 \cdot 0 + 1 + 2 \cdot 0} = -1$ (as for all SUSY partners).

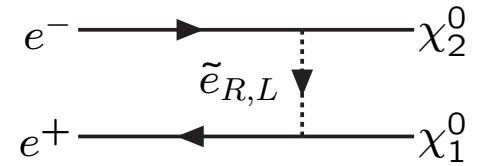
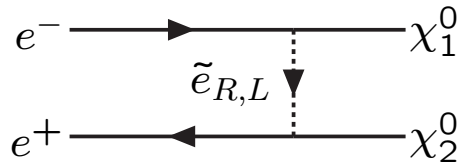
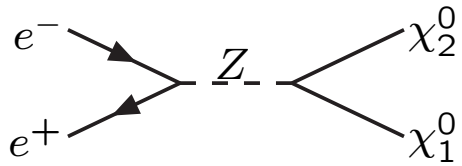
8.2 Consequences of R -parity

- proton is stable
- SUSY particles are produced in pairs
- the lightest SUSY particle (lsp) is stable
⇒ the lsp is a dark matter candidate,
the $\tilde{\chi}_1^0$ is a good candidate (uncharged, uncoloured)
- If one allows R -parity violating couplings: lepton number and baryon number violating processes are possible, bounds on R -parity violating couplings from proton decay, ...

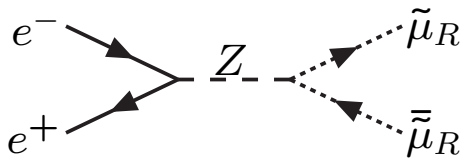
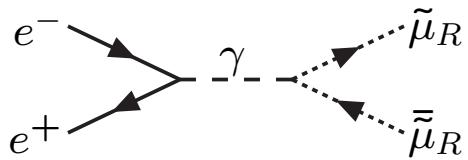
9 Examples for SUSY processes ...

At e^+e^- colliders:

Neutralino pair production:



Slepton pair production:

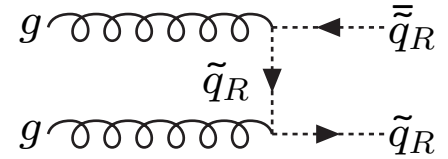
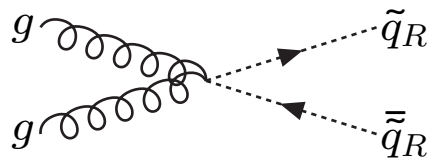
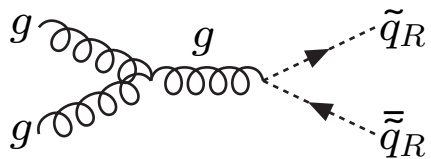
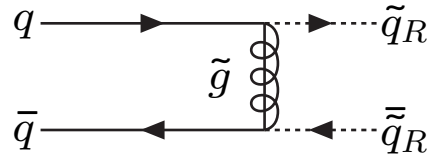
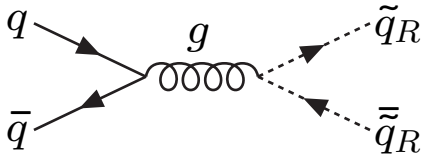


Examples for SUSY processes ...

(cont'd)

At $pp/p\bar{p}$ colliders:

Squark pair production:

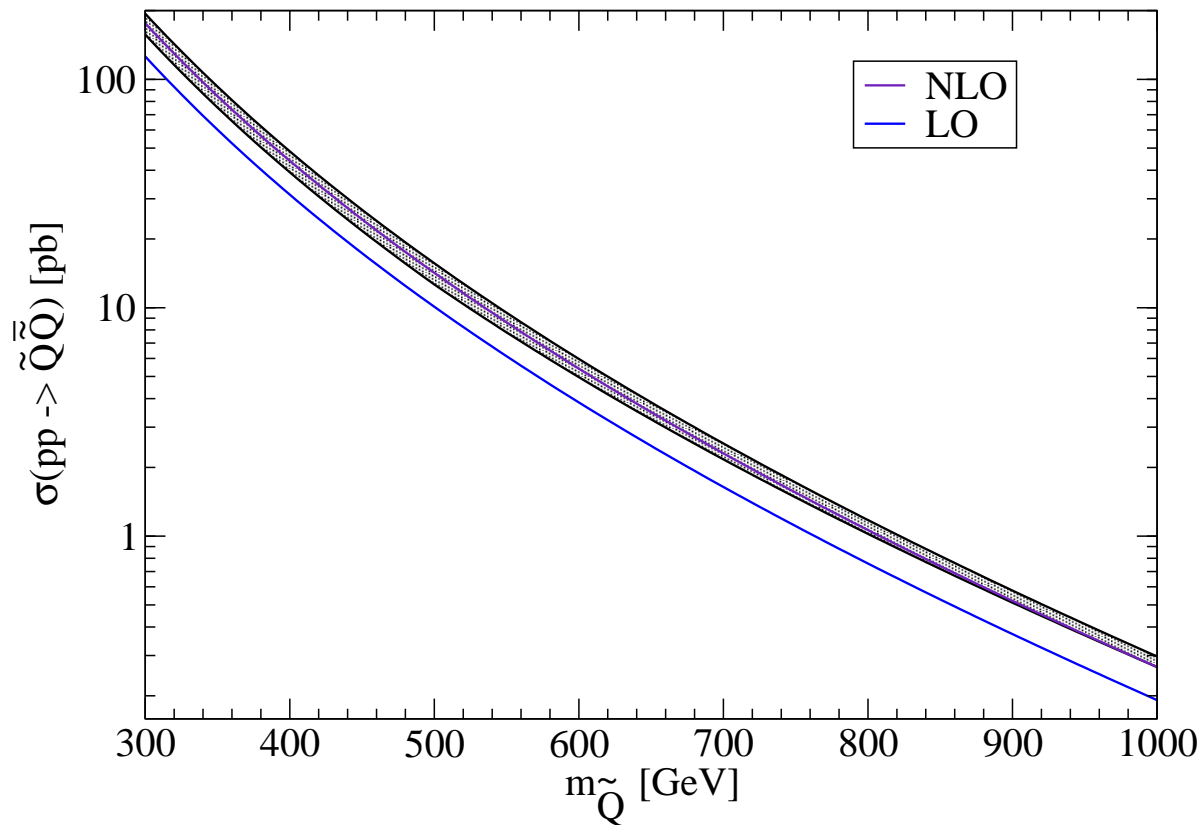


10 Cross sections for squark pair production at the LHC

$$\sigma(S, \mu^2, pp \rightarrow \tilde{Q}\bar{\tilde{Q}}) = \sum_{i,j=g,q,\bar{q}} \int_A dx_1 dx_2 p_{i,1}(x_1, \mu^2) p_{j,2}(x_2, \mu^2) \hat{\sigma}_{ij}(x_1 x_2 S, \mu^2)$$

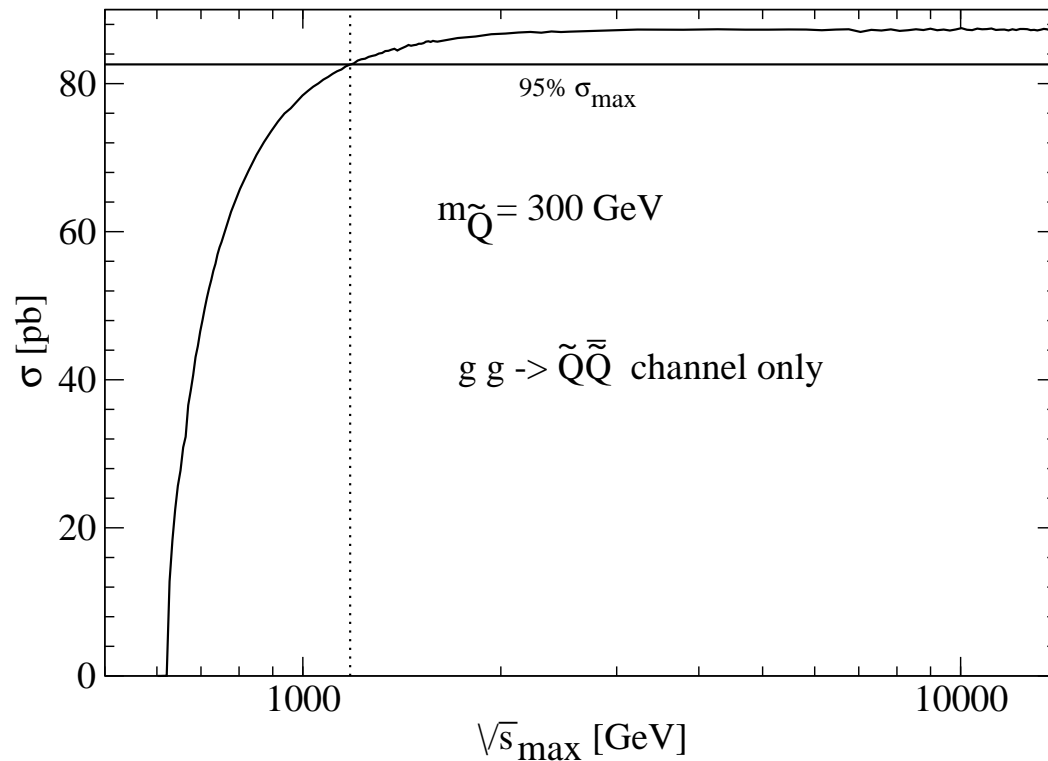
- x_i , $i = 1$ or 2 : momentum fraction of parton i
- A : Area defined by $4m^2/s \leq x_1 x_2 \leq 1$
- $\hat{\sigma}_{ij}$: partonic cross section for $ij \rightarrow \tilde{Q}\bar{\tilde{Q}}$
- $p_{i,1/2}$: parton density i , carrying momentum fraction $x_{1/2}$ of the hadron $1/2$

Cross sections for squark pair production at the LHC



Plot produced with PROSPINO [W. Beenakker, R. Höpker and M. Spira, arXiv:hep-ph/9611232.]

Cross sections for squark pair production at the LHC



for $\sqrt{s} \approx 1200$ GeV cross section saturated to 95%
 \Rightarrow Squark pairs are produced near the production threshold

Cross sections for squark pair production at the LHC (cont'd)

- Squark pairs are produced at the LHC at the production threshold with cross sections of about 1 pb...100 pb
- NLO cross sections calculated by [W. Beenakker, R. Höpker, M. Spira and P. M. Zerwas, Nucl. Phys. B **492** (1997) 51 [arXiv:hep-ph/9610490].]
- resummed NLL cross sections approximate NLO cross sections very well
want to improve cross section for squark pair production to NNLO accuracy by calculating NNLL resummed cross section → reduce scale uncertainty and improve predictive power of the cross section

11 Summary

- In the MSSM, every SM particle has its SUSY partner(s).
It is possible to include gravity, there is a dark matter candidate, and there are CP phases.
- One of the expected processes at the LHC is squark pair production
- Calculate NNLL approximated cross section