

# Structure functions and parton densities for LHC

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# Outline

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## ■ Introduction

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- Introduction
- Structure functions
  - ◆ Clean access to the structure of the proton

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- Introduction
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  - ◆ Clean access to the structure of the proton
- Parton distributions
  - ◆ NNLO evolution, sets on the market.

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- Introduction
- Structure functions
  - ◆ Clean access to the structure of the proton
- Parton distributions
  - ◆ NNLO evolution, sets on the market.
- PDFs and the LHC
  - ◆ influence on important observables,
  - ◆ more accurate extraction possible?

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- PDFs
- PDFs (2)
- Factorization Th.

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- In physics at hadron colliders the name of the game is

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**Parton Distribution Functions (PDFs)**

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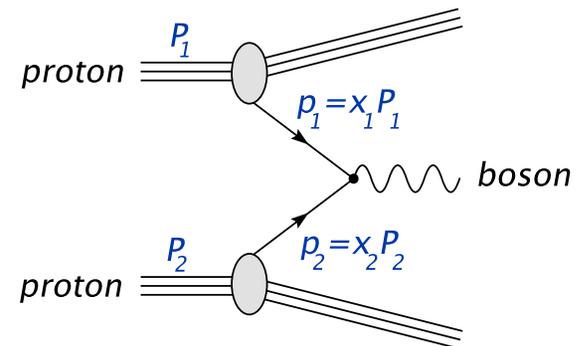
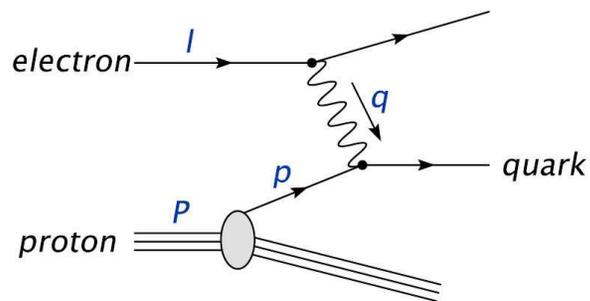
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## Parton Distribution Functions (PDFs)

- At hadron colliders we probe the non-perturbative structure of hadrons. **Examples:**



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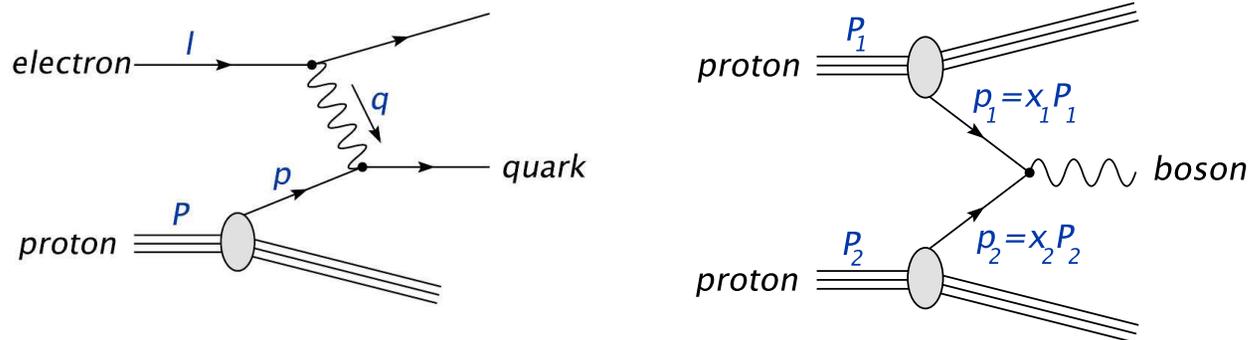
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- In physics at hadron colliders the name of the game is

## Parton Distribution Functions (PDFs)

- At hadron colliders we probe the non-perturbative structure of hadrons. **Examples:**



- All the relevant information about the structure of the particular hadron is encoded into the PDFs (details follow).

# Introduction: PDFs

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## ■ Basic properties of PDFs:

- ◆ Non-perturbative objects
- ◆ extracted from experimental data: HERA, Tevatron, etc.
- ◆ PDFs are universal: the pdf of the proton is the same in any experiment.
- ◆ This is a theoretical assumption at large enough  $Q^2$ ; consistent with experimental data.

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- Important requirement for PDFs: **precision!**

$$\Delta(O) = \Delta(\text{PDF}) + \Delta(\text{anything else})$$

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- ◆ In the past we were concerned with the mutual consistency of the PDF sets extracted under different conditions (data,  $Q^2$ ),
- ◆ The challenge of the future (**i.e. LHC**) is different: in important observables the dominant uncertainties are from PDFs (details follow).

# Introduction: Factorization Theorem

- The theoretical framework is the QCD factorization theorem

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# Introduction: Factorization Theorem

- The theoretical framework is the QCD factorization theorem
- The basic idea is that physics at very different scales factorize (i.e. does not interfere)

$$\frac{d\sigma_H}{dp_T} \simeq \sum_{i,j,k} f_i \otimes f_j \otimes \frac{d\hat{\sigma}_{i,j \rightarrow k}}{dp_T} \otimes D_{k/H} + \mathcal{O}(\Lambda_{\text{QCD}}/Q).$$

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- ◆ PDFs for each initial state hadron

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my talk tomorrow

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my talk tomorrow

- ◆ validity of the factorization approach ...

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- ◆ validity of the factorization approach ...

- Factorization Theorem is applied two ways:
  - ◆ extract PDFs from data
  - ◆ make predictions for other observables

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- SFs - measurements
- SFs - HERA
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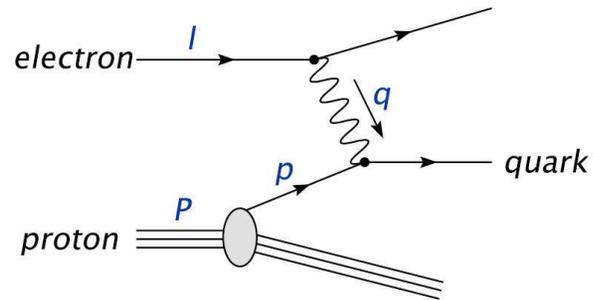
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# Structure functions

# Structure Functions

- Consider DIS in the single photon approximation:



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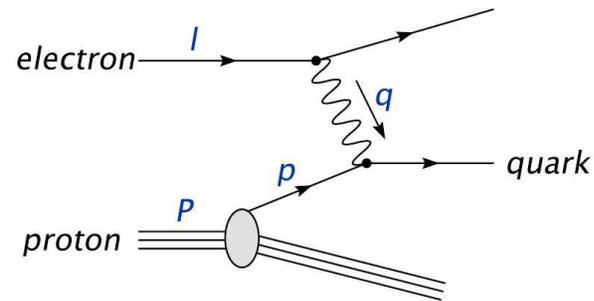
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# Structure Functions

- Consider DIS in the single photon approximation:



- The strongly interacting physics decouple from the electroweak, i.e. the cross-section takes the form:

$$d\sigma \sim L^{\mu\nu} W_{\mu\nu}$$

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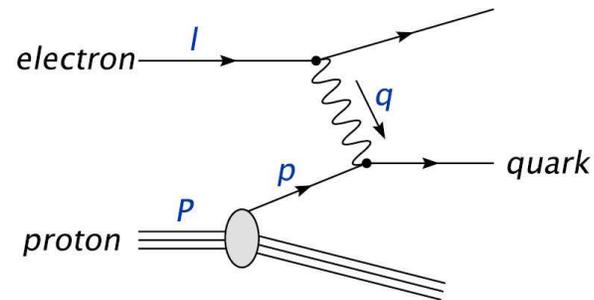
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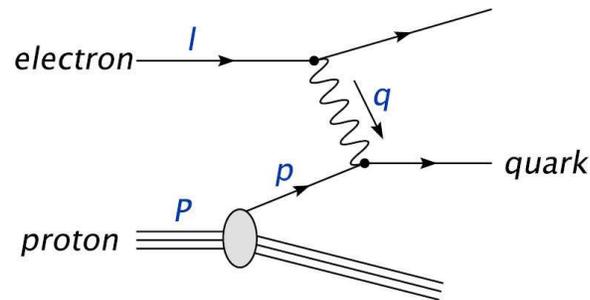
- The strongly interacting physics decouple from the electroweak, i.e. the cross-section takes the form:

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- In the unpolarized case, neglecting lepton masses parameterize  $W_{\mu\nu}$  with three structure functions  $F_1, F_2, F_3$ .

# Structure Functions

- Consider DIS in the single photon approximation:



- The strongly interacting physics decouple from the electroweak, i.e. the cross-section takes the form:

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- In the unpolarized case, neglecting lepton masses parameterize  $W_{\mu\nu}$  with three structure functions  $F_1, F_2, F_3$ .
- Both for NC and CC the measurable cross-section reads:

$$\frac{d^2\sigma}{dxdy} = \frac{4\pi\alpha^2}{xyQ^2} \left\{ y^2 x F_1 + \left[ 1 - \left( 1 + \frac{x^2 y^2 M^2}{Q^2} \right) \right] F_2 \pm xy \left( 1 - \frac{y}{2} \right) F_3 \right\}$$

# Structure Functions: measurements

## ■ What can be measured and to what extend?

Courtesy PDG and H1

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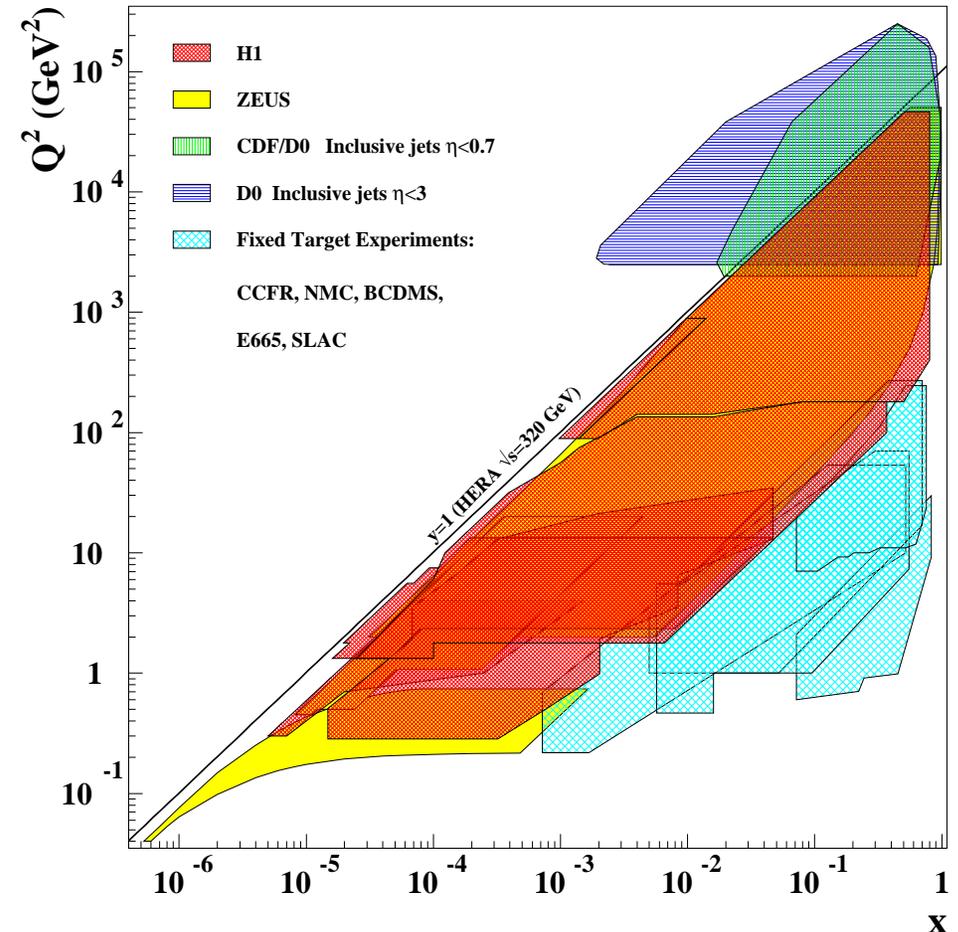
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Process	Main Subprocess	PDFs Probed
$\ell^\pm N \rightarrow \ell^\pm X$	$\gamma^* q \rightarrow q$	$g(x \lesssim 0.01), q, \bar{q}$
$\ell^+ (\ell^-) N \rightarrow \bar{\nu} (\nu) X$	$W^* q \rightarrow q'$	
$\nu (\bar{\nu}) N \rightarrow \ell^- (\ell^+) X$	$W^* q \rightarrow q'$	
$\nu N \rightarrow \mu^+ \mu^- X$	$W^* s \rightarrow c \rightarrow \mu^+$	$s$
$\ell N \rightarrow \ell Q X$	$\gamma^* Q \rightarrow Q$	$Q = c, b$
	$\gamma^* g \rightarrow Q\bar{Q}$	$g(x \lesssim 0.01)$
$pp \rightarrow \gamma X$	$q\bar{q} \rightarrow \gamma q$	$g$
$pN \rightarrow \mu^+ \mu^- X$	$q\bar{q} \rightarrow \gamma^*$	$\bar{q}$
$pp, pn \rightarrow \mu^+ \mu^- X$	$u\bar{u}, d\bar{d} \rightarrow \gamma^*$	$\bar{u} - \bar{d}$
	$u\bar{d}, d\bar{u} \rightarrow \gamma^*$	
$ep, en \rightarrow e\pi X$	$\gamma^* q \rightarrow q$	
$p\bar{p} \rightarrow W \rightarrow \ell^\pm X$	$ud \rightarrow W$	$u, d, u/d$
$p\bar{p} \rightarrow \text{jet} + X$	$gg, qg, qq \rightarrow 2j$	$q, g(0.01 \lesssim x \lesssim 0.5)$



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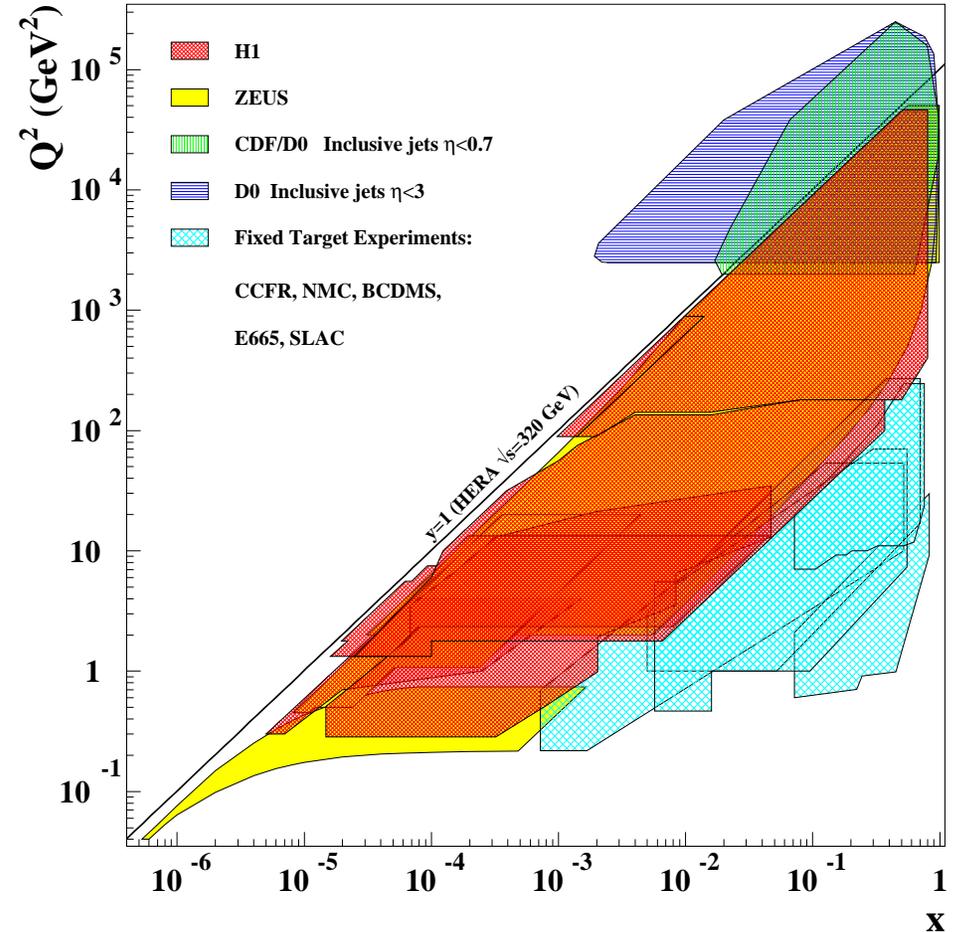
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$p\bar{p} \rightarrow \text{jet} + X$	$gg, qg, qq \rightarrow 2j$	$q, g(0.01 \lesssim x \lesssim 0.5)$



## ■ Possible to test DGLAP evolution across a range of scales and experiments!

# Structure Functions: results from HERA

- HERA (H1 and ZEUS) provide us with a wealth of data on both NC and CC DIS.

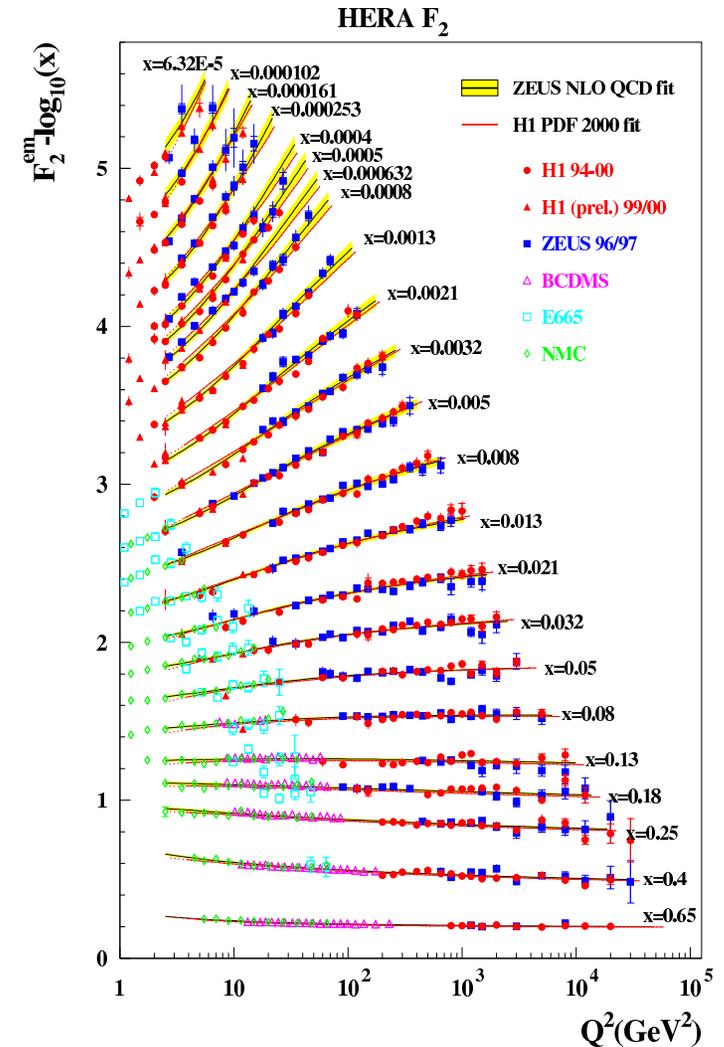
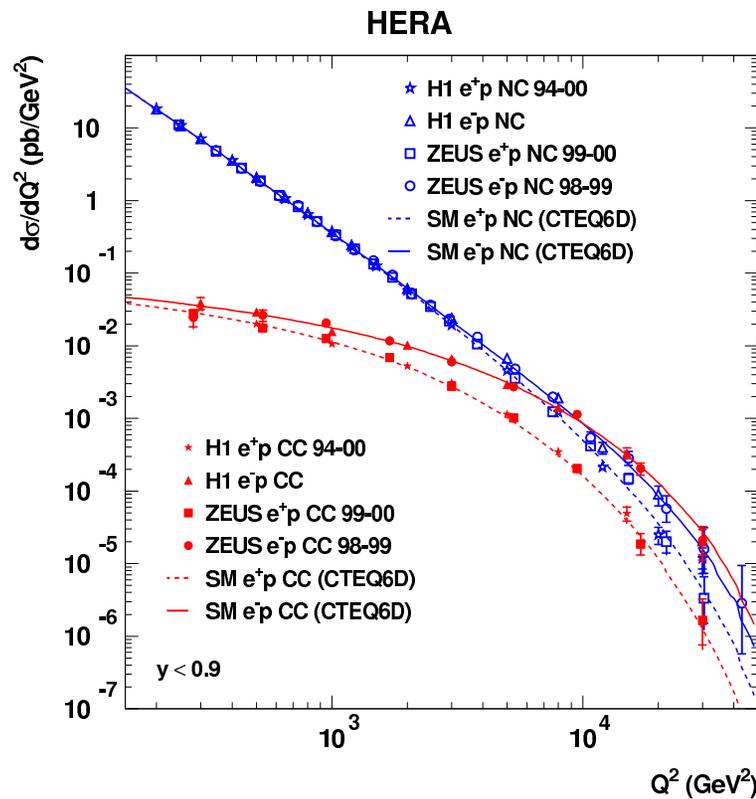
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# Structure Functions: fixed target; NuTeV

- Important results from fixed target experiments.

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  - ◆  $\nu$  and  $\bar{\nu}$  -  $N$  experiment,

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  - ◆  $10^5$   $\bar{\nu}$  -  $N$  and  $10^6$   $\nu$  -  $N$  CC events,

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  - ◆  $10^5$   $\bar{\nu}$  -  $N$  and  $10^6$   $\nu$  -  $N$  CC events,
  - ◆ measured NC/CC ratio:

$$R^- = \frac{\sigma_{NC}^\nu - \sigma_{NC}^{\bar{\nu}}}{\sigma_{CC}^\nu - \sigma_{CC}^{\bar{\nu}}} \sim \frac{1}{2} - \sin^2 \theta_W$$

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- ◆ the measurement was in  $3\sigma$  difference with the "SM".

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- Possible explanations?

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- ◆ the measurement was in  $3\sigma$  difference with the "SM".
- Possible explanations?
  - ◆ BSM physics,

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# Structure Functions: fixed target; NuTeV

- Important results from fixed target experiments.

- NuTeV:

- ◆  $\nu$  and  $\bar{\nu}$  -  $N$  experiment,
- ◆ Iron target - nuclear corrections.
- ◆  $10^5$   $\bar{\nu} - N$  and  $10^6$   $\nu - N$  CC events,
- ◆ measured NC/CC ratio:

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Bardin, Dokuchaeva ('86); Diener, Dittmaier, Hollik ('03,'05)

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Bardin, Dokuchaeva ('86); Diener, Dittmaier, Hollik ('03,'05)

- ◆ strange sea asymmetry:  $s - \bar{s} \neq 0$

Brodsky, Ma ('96); Catani, de Florian, Rodrigo, Vogelsang ('04)

# Structure Functions: QCD theory

- What does QCD say about the Structure Functions?

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# Structure Functions: QCD theory

- What does QCD say about the Structure Functions?
- Actually - a lot. From the factorization theorem:

$$F_n^{(NC,CC)}(x, Q^2) = \sum_{i=q,g} f_i(x, Q^2) \otimes C_{i,n}^{(NC,CC)}(x, Q^2)$$

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- That factorization requires a choice of scheme:
  - ◆  $\overline{\text{MS}}$ ,
  - ◆ **DIS**

# Structure Functions: QCD theory (cont.)

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- What is currently known in DIS?

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## ■ What is currently known in DIS?

- ◆ massless inclusive NC DIS: 3 loop (NNNLO)

Moch, Vermaseren, Vogt ('05)

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## ■ What is currently known in DIS?

- ◆ massless inclusive NC DIS: 3 loop (NNNLO)
- ◆ massless inclusive CC DIS: 2 loop (NNLO)

Moch, Vermaseren, Vogt ('05)

Zijlstra, van Neerven ('92)

# Structure Functions: QCD theory (cont.)

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◆ massless inclusive NC DIS: 3 loop (NNNLO)

Moch, Vermaseren, Vogt ('05)

◆ massless inclusive CC DIS: 2 loop (NNLO)

Zijlstra, van Neerven ('92)

◆ massive inclusive NC and CC DIS: 1 loop (NLO)

Glück, Kretzer, Reya ('96); Kretzer, Schienbein ('98)

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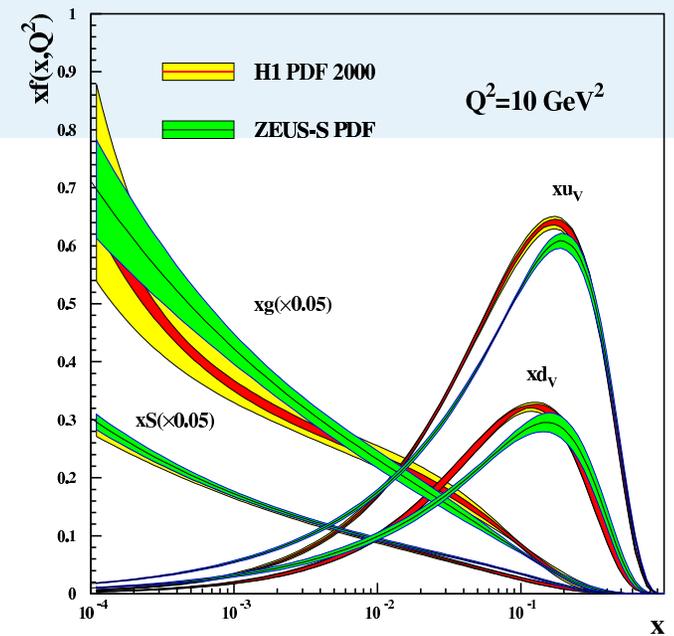
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# Parton Distribution Functions

# DGLAP equation

- PDFs are purely non-perturbative



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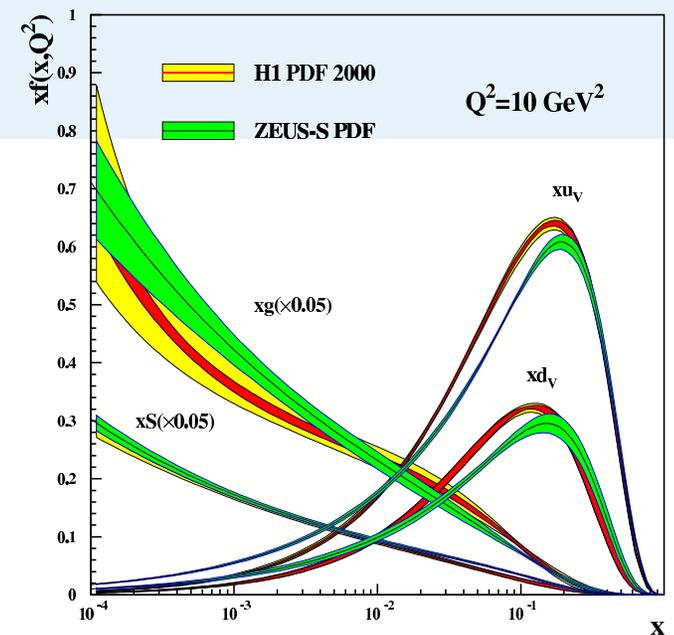
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# DGLAP equation

■ PDFs are purely non-perturbative

■ Their scale evolution is perturbative!  
DGLAP evolution equation:

$$\frac{df_i(x, \mu^2)}{d \ln(\mu^2)} = \sum_j P_{ij}^S(\alpha_s(\mu^2), x) \otimes f_j(x, \mu^2) ; \quad i, j = q, \bar{q}, g.$$



Altarelli, Parisi ('77); Gribov, Lipatov ('72); Dokshitzer ('77)

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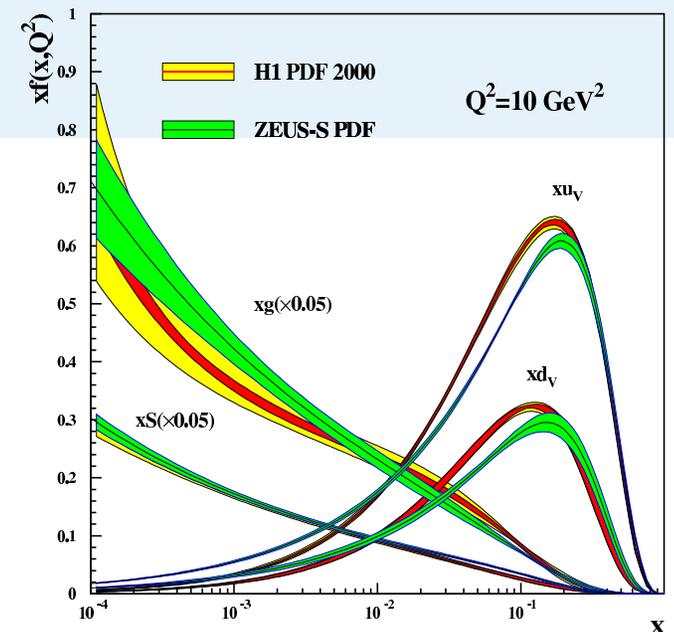
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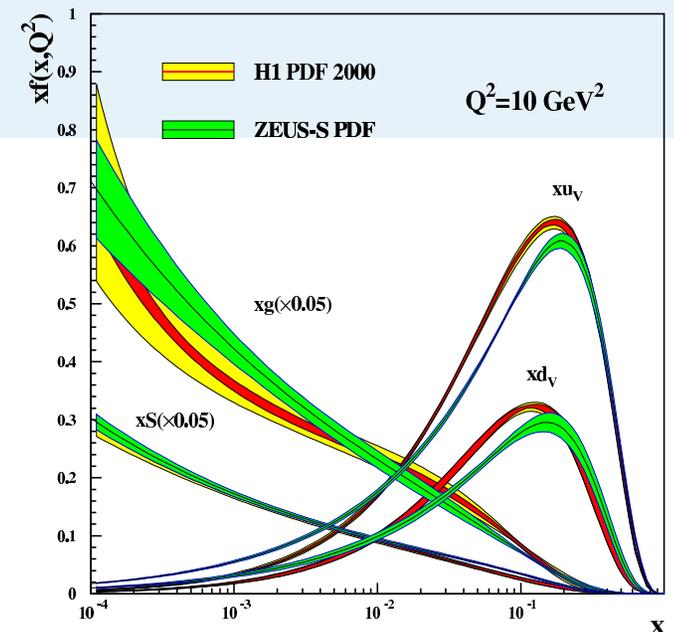
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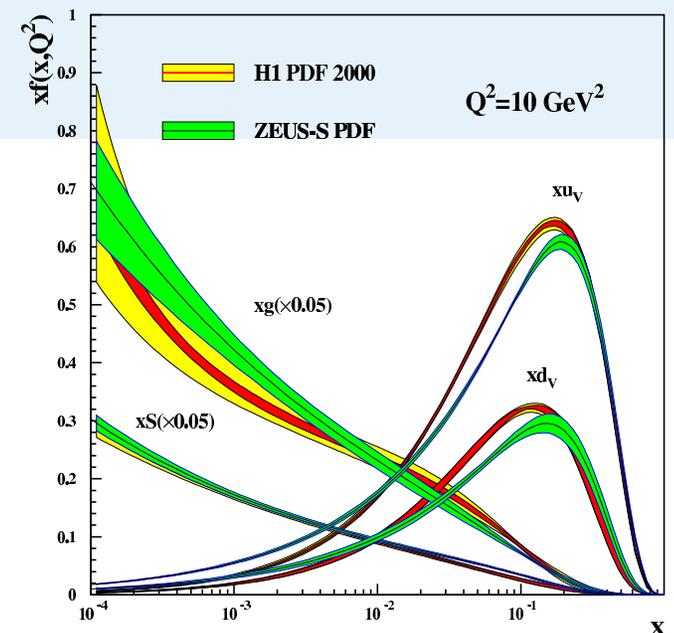
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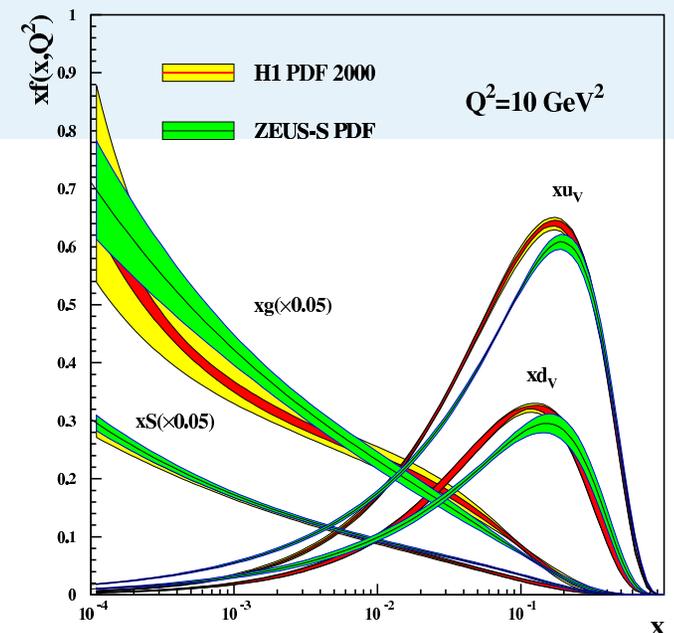
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Altarelli, Parisi (77); Gribov, Lipatov (72); Dokshitzer (77)

- DGLAP is a large system of integro-differential equations.
- Ways to solve it:
  - ◆ partially diagonalize it by splitting in singlets and non-singlets,
  - ◆ In Mellin space the integrals convolution becomes ordinary multiplication:

$$f \otimes g(x) \rightarrow f \cdot g(N)$$

# Solving DGLAP

- Solve analytically in Mellin space, then invert numerically to  $x$ -space.

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# Solving DGLAP

- Solve analytically in Mellin space, then invert numerically to  $x$ -space.
- Publicly available ([here](#)) evolution code for parton densities to NNLO: **QCD-Pegasus**.

A. Vogt ('04)

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- To completely specify the solution, an initial condition at some scale is required.
- The initial scale is usually chosen low, around  $1 - 2\text{GeV}$ .
- Examples:
  - ◆ MRST:  $\mu_0^2 = 1\text{GeV}^2$   
 $xg(x) = c_1 x^a (1-x)^b (1 + c_3 x^{0.5} + c_4 x) - c_5 x^c (1-x)^d$ .
  - ◆ CTEQ6:  $\mu_0^2 = 1.69\text{GeV}^2$   
 $xg(x) = c_1 x^a (1-x)^b e^{cx} (1 + c_2 x)^d$ .

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# Splitting functions of QCD

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DGLAP equation is specified by the QCD splitting functions:

$$P_{ij}^{(S,T)}(\alpha_s(\mu^2), x) = \sum_{n=0} \left( \frac{\alpha_s(\mu^2)}{2\pi} \right)^{n+1} P_{ij}^{(n)(S,T)}(x)$$

Beyond LO, the Time- and Space-like functions differ!

■ Space-like kinematics  
(DIS)

■ Time-like kinematics  
( $e^+e^-$ )

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## ■ Space-like kinematics (DIS)

### ■ LO

Altarelli, Parisi ('77)

### ■ NLO

Curci, Furmanski, Petronzio ('80)

## ■ Time-like kinematics ( $e^+e^-$ )

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Altarelli, Parisi ('77)

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## ■ Space-like kinematics (DIS)

### ■ LO

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### ■ NLO

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Moch, Vermaseren, Vogt ('04)

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### ■ LO

Altarelli, Parisi ('77)

### ■ NLO

Curci, Furmanski, Petronzio ('80)

### ■ NNLO - NS available

A.M., Moch, Vogt ('06)

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# Splitting functions of QCD (cont.)

## ■ Calculation of the splitting functions

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- Calculation of the splitting functions
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# Splitting functions of QCD (cont.)

- Calculation of the splitting functions
  - ◆ Take any collinear observable (like inclusive DIS),
  - ◆ regularize all soft/collinear singularities dimensionally,
  - ◆ the  $n$ -loop splitting functions are the coefficients of

$$\sim \frac{1}{\epsilon} \alpha_s^n$$

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# Splitting functions of QCD (cont.)

- Calculation of the splitting functions
  - ◆ Take any collinear observable (like inclusive DIS),
  - ◆ regularize all soft/collinear singularities dimensionally,
  - ◆ the  $n$ -loop splitting functions are the coefficients of

$$\sim \frac{1}{\epsilon} \alpha_s^n$$

- In the calculation of the space-like functions  
Moch, Vermaseren and Vogt used

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- In the calculation of the space-like functions Moch, Vermaseren and Vogt used
  - ◆ inclusive DIS
  - ◆ OPE and worked directly in Mellin moment space
- To get the time-like NS functions a relation between the NS time-like and space-like functions was devised.

A.M., Moch, Vogt ('06)

The result can be cast in a **very** simple form:

Dokshitzer, Marchesini, Salam ('05)

$$P^{(2),T} - P^{(2),S} = \left[ \ln(z) \cdot \left( P^{(1),T} + P^{(1),S} \right) \right] \otimes P^{(0)} + \left[ P^{(1),T} + P^{(1),S} \right] \otimes \left[ \ln(z) \cdot P^{(0)} \right]$$

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- In principle the extracted PDFs should not depend on the details of the extraction.

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# Sets of PDFs

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- However, they do:

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# Sets of PDFs

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  - ◆ **GRV**
  - ◆ **ZEUS**
- Typically, the error bands are obtained by varying the parameters in the parametric form at  $\mu_0^2$  ...

# Sets of PDFs (cont.)

- A great website with online plotting and comparison between PDF sets.

<http://durpdg.dur.ac.uk/hepdata/pdf3.html>

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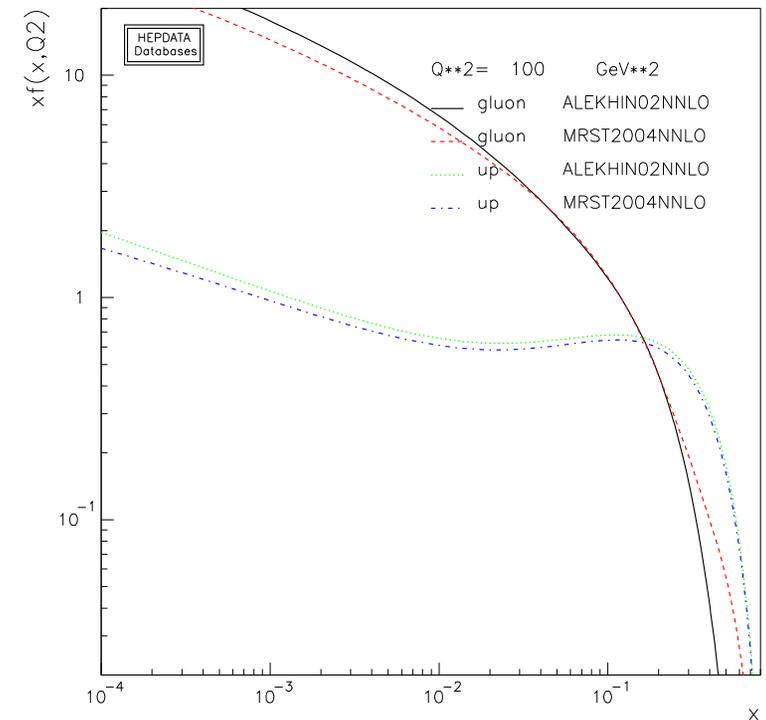
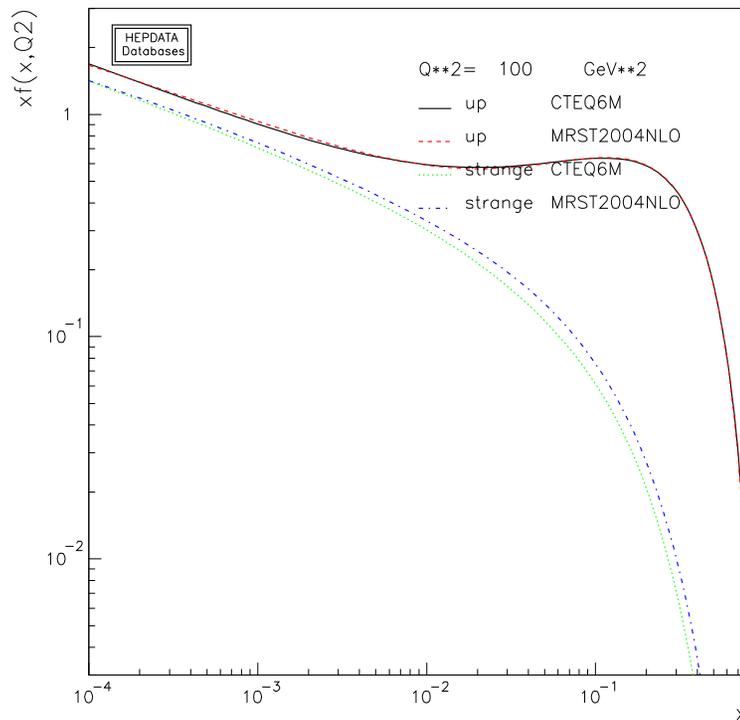
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# Sets of PDFs (cont.)

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- Error plotting is optional.



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# Sets of PDFs: comparisons, error bands

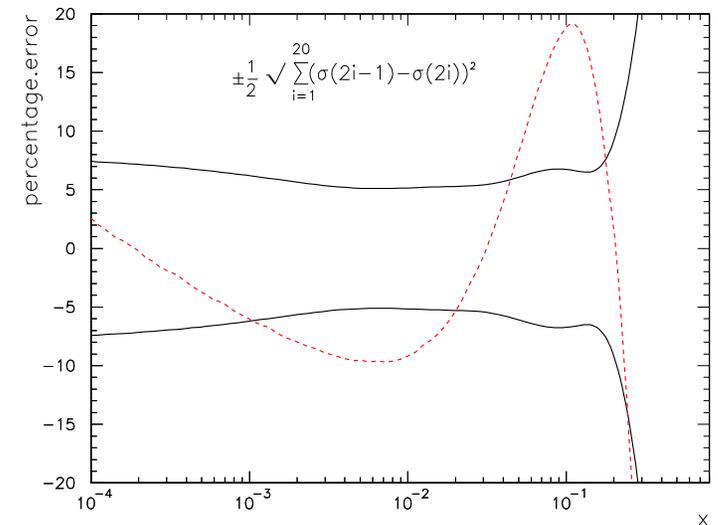
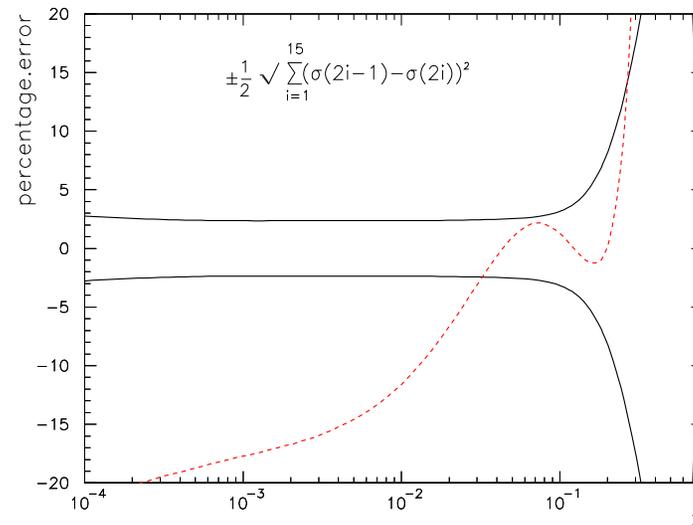
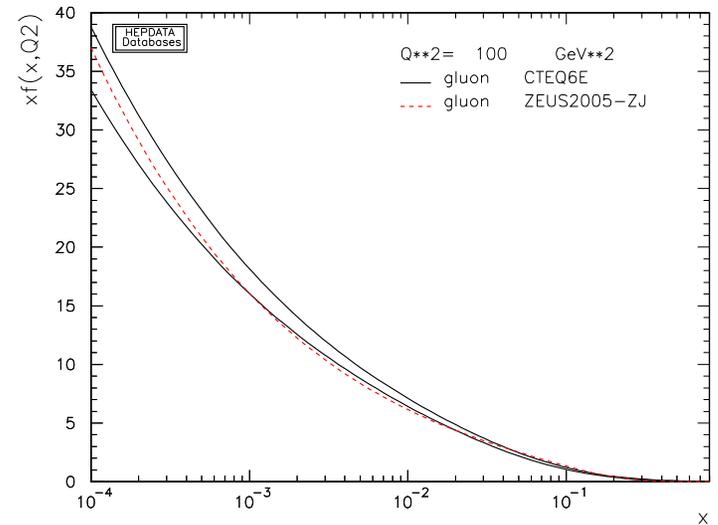
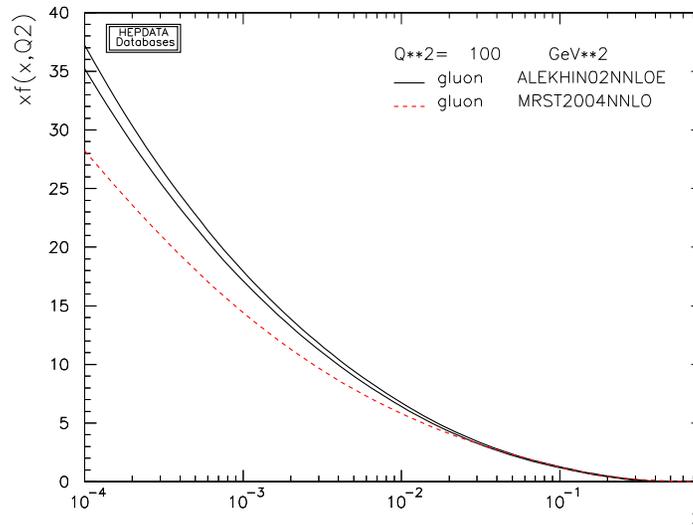
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# Sets of PDFs: comparisons, errors (cont.)

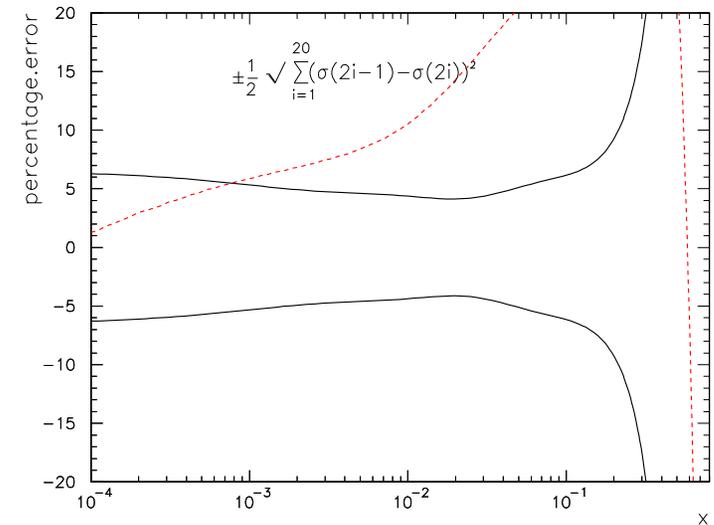
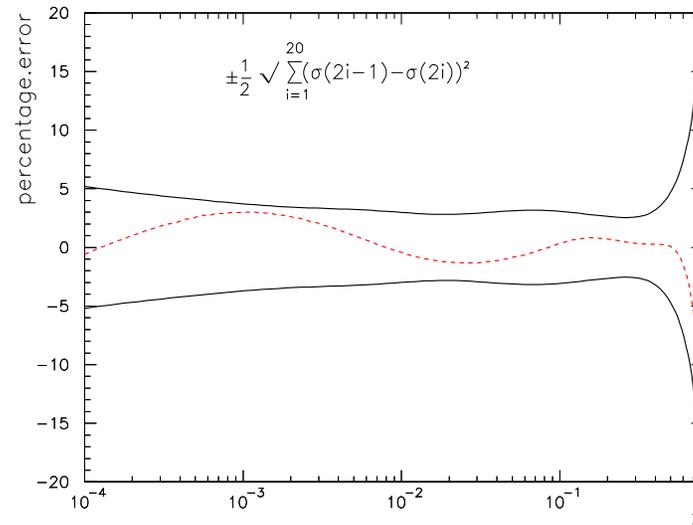
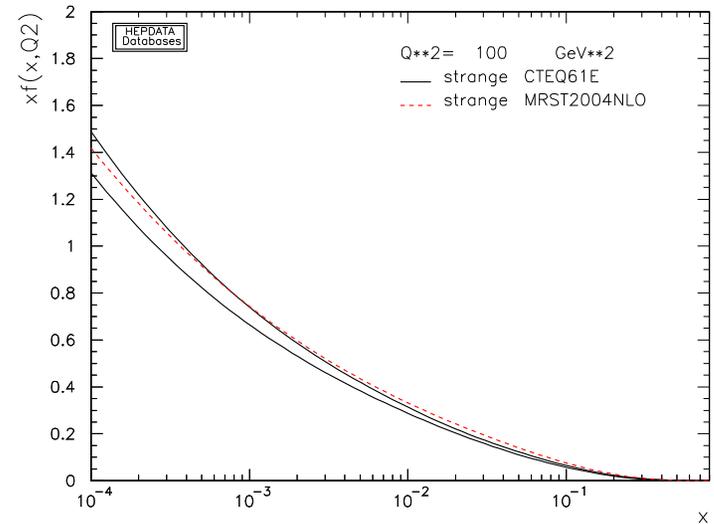
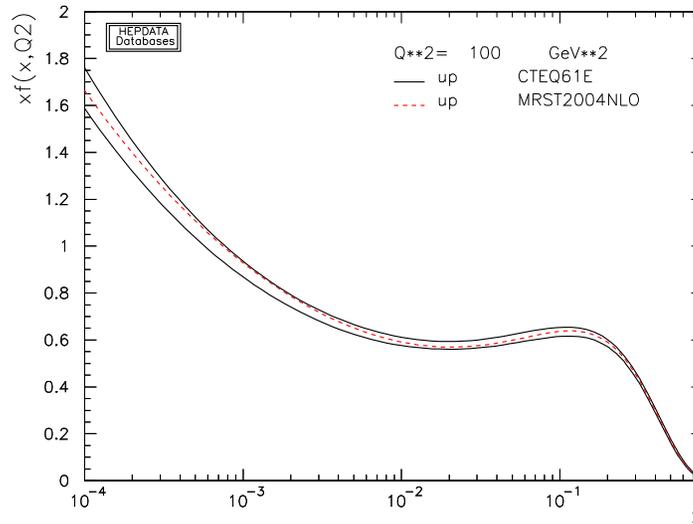
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**PDFs and the LHC**

- PDFs for LHC: W/Z/DY
- PDFs and LHC: Higgs
- PDFs and LHC: other

# PDFs and the LHC

# PDFs for LHC: W/Z/Drell-Yan

- We need precision for the LHC because they directly influence important observables there:

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PDFs and the LHC

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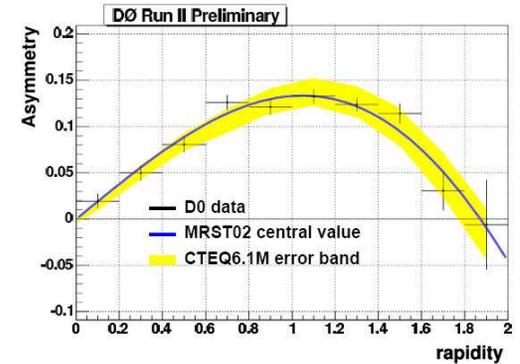
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M. Sanders, talk at ICHEP2006

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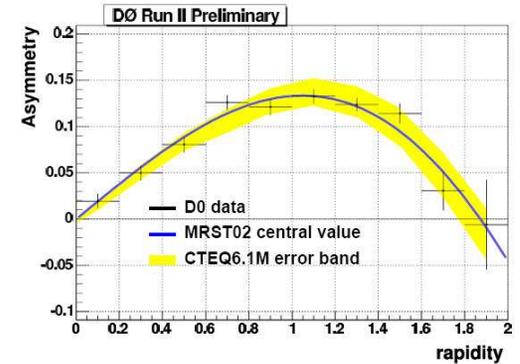
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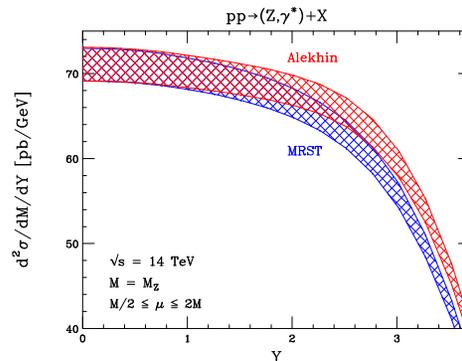
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- ◆ The very small perturbative uncertainty makes it possible to distinguish current PDFs (plots for LHC):

Anastasiou, Dixon, Melnikov, Petriello ('03)

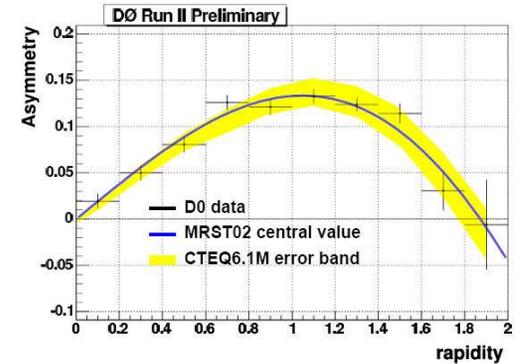


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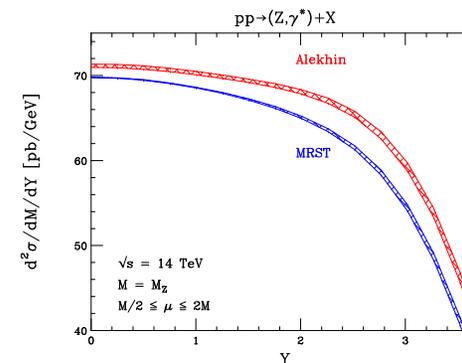
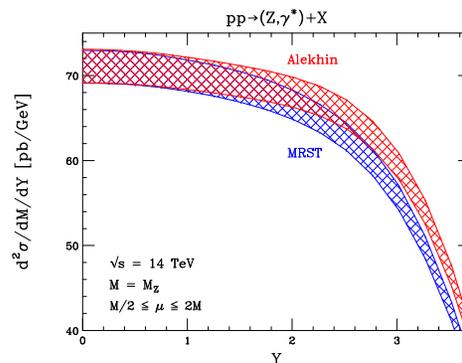
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# PDFs and LHC: Higgs

- Higgs production - couples (through a top loop) to the gluon

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● PDFs for LHC: W/Z/DY

● PDFs and LHC: Higgs

● PDFs and LHC: other

# PDFs and LHC: Higgs

- Higgs production - couples (through a top loop) to the gluon
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Harlander, Kilgore ('02); Anastasiou, Melnikov ('02); Ravindran, Smith, van Neerven ('03)

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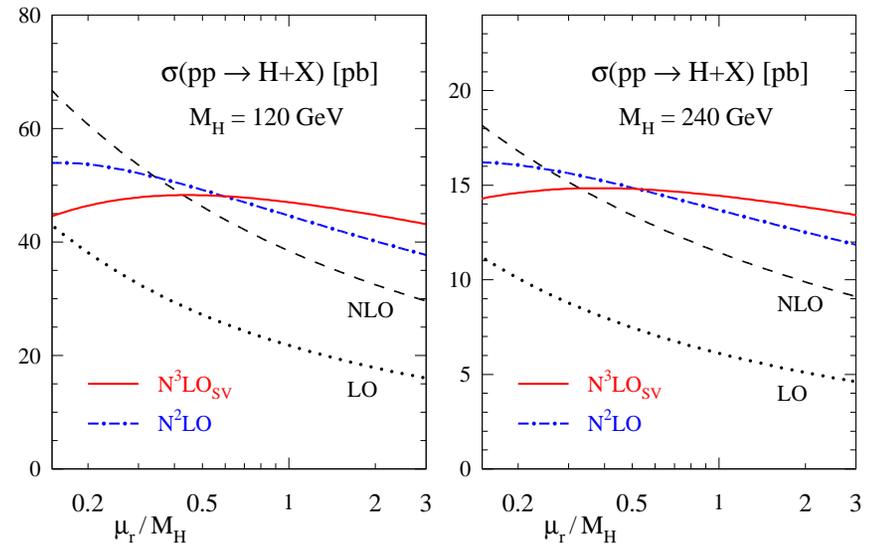
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- and  $N^3\text{LO}$ -(soft+virtual) corrections known

Moch, Vogt ('05)



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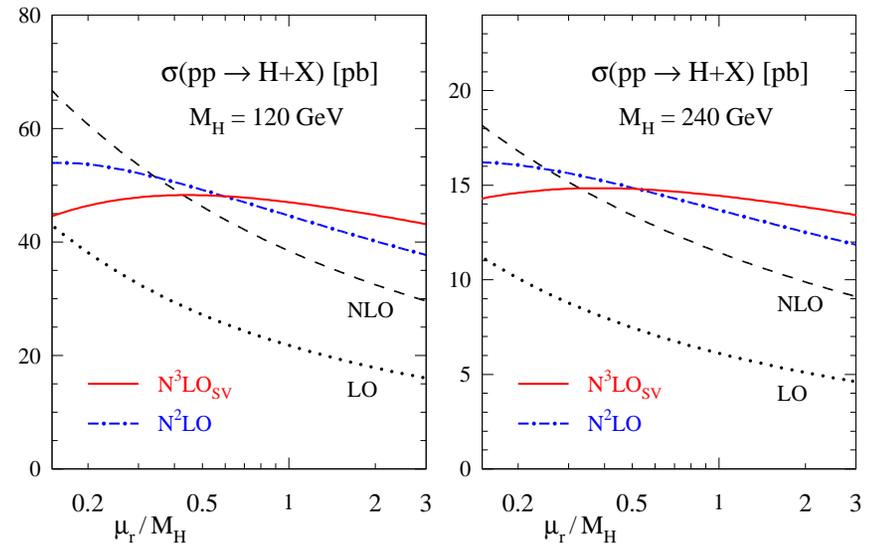
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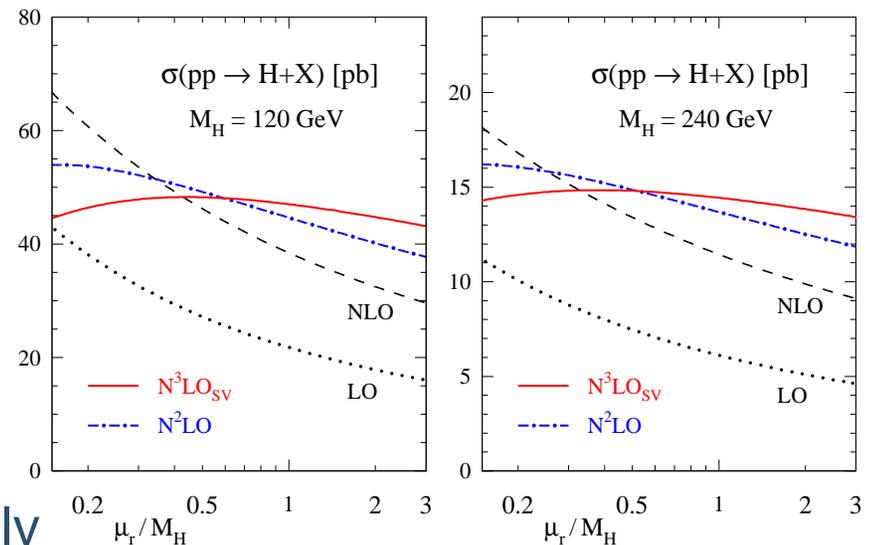
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- Moreover, the decay  $pp \rightarrow H + X \rightarrow \gamma\gamma$  is now known fully differentially



Anastasiou, Melnikov, Petriello ('05)

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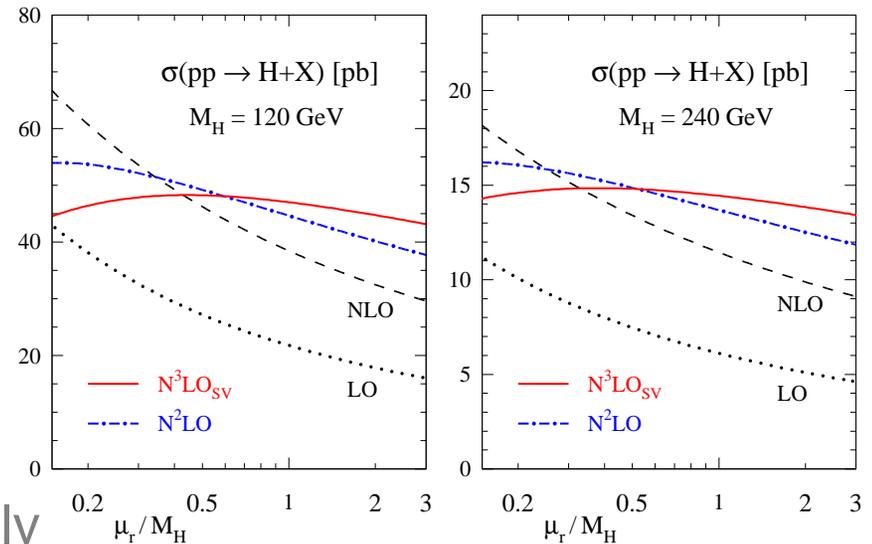
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Anastasiou, Melnikov, Petriello ('05)

- Perturbative uncertainty in the 8 – 10%. Perhaps even less.

Catani, de Florian, Grazzini, Nason ('03)

# PDFs and LHC: Higgs, $t\bar{t}$ , jets.

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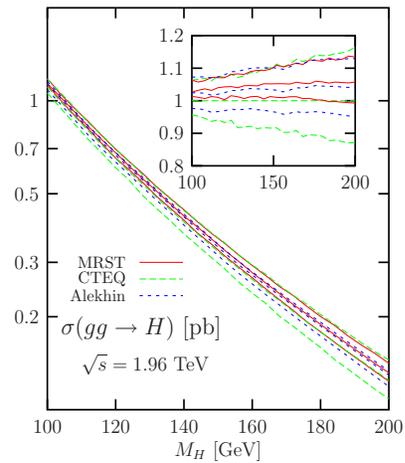
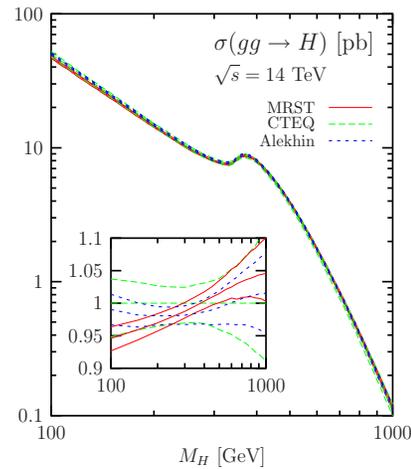
PDFs and the LHC

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● PDFs and LHC: Higgs

● PDFs and LHC: other

- Uncertainties from PDF are 3 – 5%. However, set dependent:



Djouadi, Ferrag ('03)

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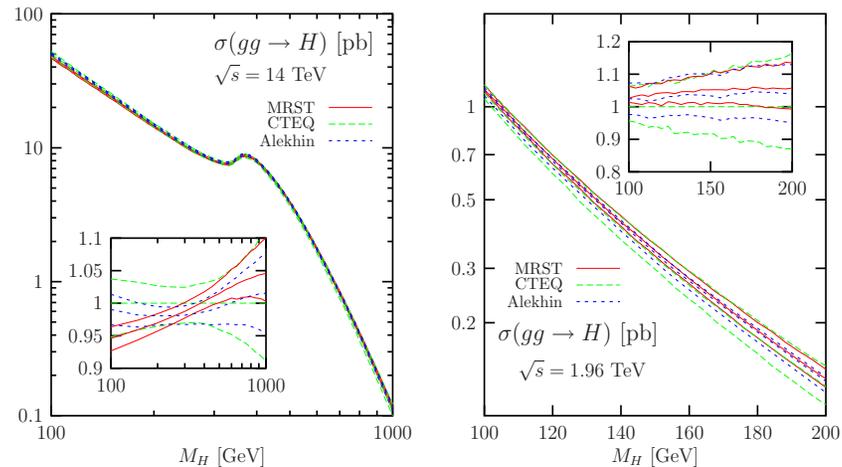
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Djouadi, Ferrag ('03)

- The gluon distribution is important for  $t\bar{t}$  too.
  - ◆ PDF uncertainty  $\sim 3 - 4\%$ ,
  - ◆ NLO + NLL  $\sim 10\%$ ?
  - ◆ total experimental  $\lesssim 10\%$ .

talk by Peter Uwer

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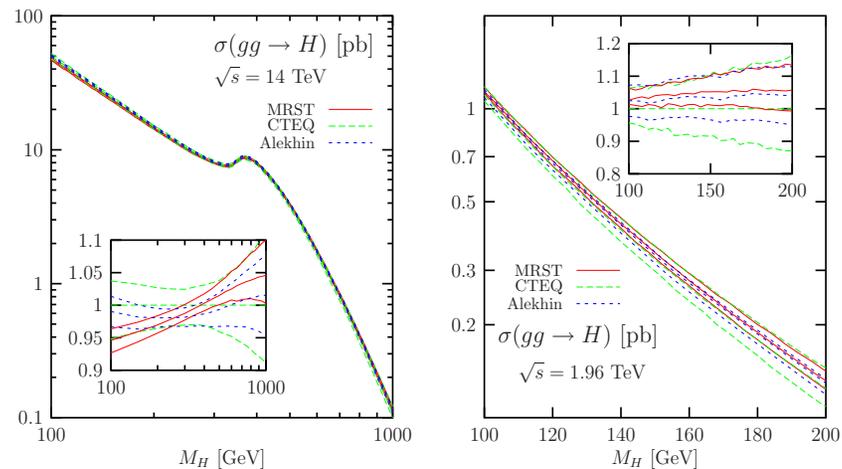
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  - ◆ total experimental  $\lesssim 10\%$ .

talk by Peter Uwer

- jet production at NNLO will be important to constrain it better at large  $x$ .

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● PDFs for LHC: W/Z/DY

● PDFs and LHC: Higgs

● PDFs and LHC: other

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talk by Jochen Bartels today ...

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  - ◆  $b, c$ .
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  - ◆ ... it is feasible!
- LHC is behind the corner. A great challenge for all of us!

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