

Where is the PdV term in the first law of black hole thermodynamics?

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Outline

Review of black hole thermodynamics

First Law

Smarr Relation

Enthalpy and Volume

Enthalpy

Equation of State - Van der Waals

Compressibility

de Sitter

Summary

Black hole thermodynamics

- Entropy, $S = \frac{1}{4} \frac{A}{\ell_{Pl}^{D-2}}$: $A = \text{area}$, ($\ell_{Pl}^{D-2} = G_N \hbar$, $c = 1$).
- Hawking temperature, $T = \frac{\kappa \hbar}{2\pi}$, $\kappa = \text{surface gravity}$.
- Internal (thermal) energy, identify $M = U(S)$,

First Law of Black Hole Thermodynamics

$$dM = dU = T dS = \frac{\kappa}{8\pi G_N} dA.$$

- Including J and Q ,

$$dM = dU = T dS + \Omega dJ + \Phi dQ$$

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Smarr relation

- Ordinary thermodynamics: $U(S, V, N)$ (N = no. of moles) is a function of **extensive variables**. U is also extensive.

$$\lambda^d U(S, V, N) = U(\lambda^d S, \lambda^d V, \lambda^d N)$$

$$\Rightarrow U = S \frac{\partial U}{\partial S} + V \frac{\partial U}{\partial V} + N \frac{\partial U}{\partial N} \quad \text{Euler equation}$$

$$\Rightarrow U = ST - VP + N\mu \quad (\mu = \text{chemical potential})$$

$$\Rightarrow G = N\mu = U - ST + VP \quad \text{Gibbs-Duhem relation}$$

Black hole, ADM mass M :

$$S \rightarrow \lambda^{D-2} S, J \rightarrow \lambda^{D-2} J, M \rightarrow \lambda^{D-3} M,$$

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Doesn't work when $\Lambda \neq 0$!

(Smarr 1973)

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First Law

BPD [1008.5023].

$$dU = T dS + \Omega dJ - P dV$$

- $M = U + PV$ with $V = \frac{\partial M}{\partial P}$, thermodynamic volume.

e.g. for asymptotically AdS, $D = 4$:

$\eta_{max} = 0.5184 \dots$, for extremal black hole.

$\eta_{max} = 1 - \frac{1}{\sqrt{2}} = 0.2929 \dots$, without the PdV term.

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Thermodynamic volume

- Asymptotically anti-de Sitter Kerr space-time,
 $\Lambda < 0 \Rightarrow P > 0$.
- black hole event horizon, r_h ;
“geometric volume”, $V_0 = \frac{r_h A_h}{3}$. (D=4)

Thermodynamic volume

$$V = V_0 + \frac{4\pi}{3} \frac{J^2}{M} > V_0.$$

- Spherically symmetric case: $J = 0$, $V = \frac{4\pi r_h^3}{3}$.
- Reverse isoperimetric inequality
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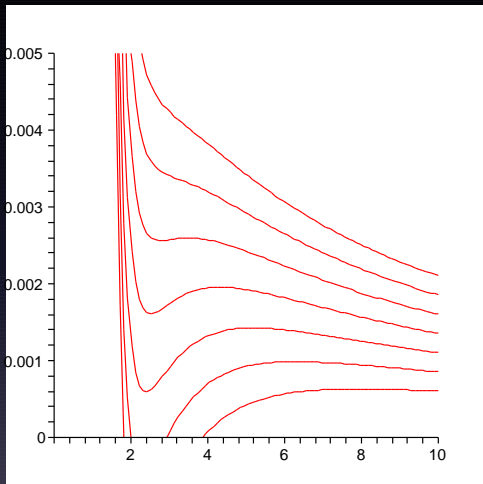
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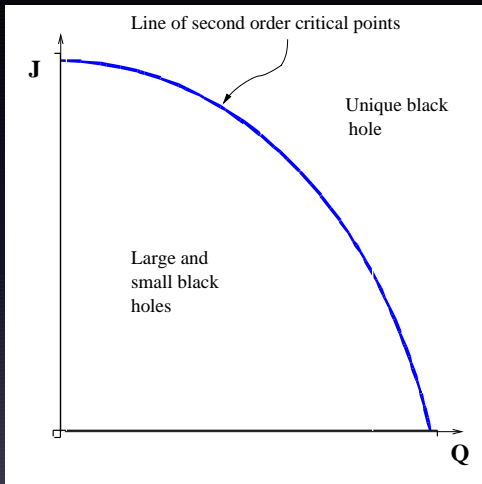
Equation of state: P–V diagram



P as a function of $\left(\frac{3V}{4\pi}\right)^{1/3}$, curves of constant T for $J = 1$.

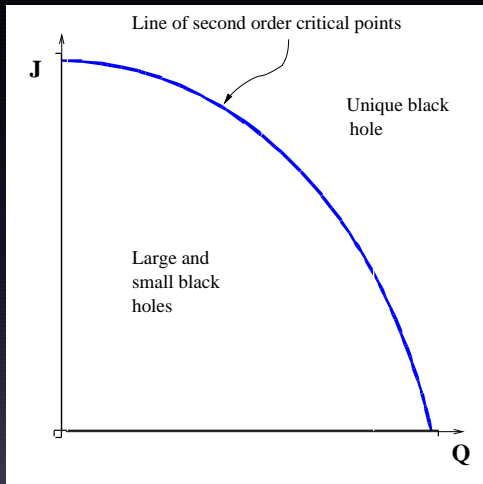
Critical point at $T_C = 0.0413/J^{1/2}$, $P_C = 0.00280/J$ and $V_C = 12.90J^{3/2}$, Caldarelli, Gognola+Klemm (1999).

Kerr-Reissner-Nordström-AdS



- Reissner-Nordström anti-de Sitter ($J \neq 0$, $Q \neq 0$)
- Mean field exponents — same as Van der Waals gas, Emparan, Johnson, and Myers (1999); Gunasekaran, Kubizňák and Mann; BPD (2012).

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Compressibility

BPD [arXiv:1109.0198]

- Adiabatic compressibility: $\kappa = -\frac{1}{V} \left. \frac{\partial V}{\partial P} \right|_{S,J} \geq 0$.

- $\kappa_{J=0} = 0$.

- Maximum for J_{max} ($T = 0$): $\kappa_{max} = \frac{2S(1+8PS)}{(3+8PS)^2(1+4PS)}$.

- e.g. $P = 0$,

$$\kappa_{max} = \frac{2S}{9} = \frac{4\pi M^2}{9} = 2.6 \times 10^{-38} \left(\frac{M}{M_{\odot}} \right)^2 m s^2 kg^{-1}.$$

cf. neutron star, $M \approx M_{\odot}$, $R \approx 10 km$, degenerate Fermi gas $\Rightarrow \kappa \approx 10^{-34} m s^2 kg^{-1}$.

Very stiff equation of state!

Asymptotically de Sitter

BPD, D. Kastor, D. Kubiznak, R.B. Mann and J. Traschen [1301.5926]

- $P = -\frac{\Lambda}{8\pi} < 0$.
- Two event horizons: black hole r_h ; cosmological, r_c .
- Two different temperatures, $T_h \neq T_c$, in general.
- $M(S_h, P, J) = M(S_c, P, J)$.

$$V_h = \left. \frac{\partial M}{\partial P} \right|_{S_h, J}, \quad V_c = \left. \frac{\partial M}{\partial P} \right|_{S_c, J}.$$

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For fixed V_c the cosmological horizon entropy, $S_c = \frac{A_c}{4}$, is maximized by Schwarzschild- de Sitter space-time ($J = 0$).

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For fixed V_c the cosmological horizon entropy, $S_c = \frac{A_c}{4}$, is maximized by Schwarzschild- de Sitter space-time ($J = 0$).

- Volume between horizons: $V = V_c - V_h = \frac{r_c A_c}{3} - \frac{r_h A_h}{3}$.
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Asymptotically de Sitter

BPD, D. Kastor, D. Kubiznak, R.B. Mann and J. Traschen [1301.5926]

- $P = -\frac{\Lambda}{8\pi} < 0$.
- Two event horizons: black hole r_h ; cosmological, r_c .
- Two different temperatures, $T_h \neq T_c$, in general.
- $M(S_h, P, J) = M(S_c, P, J)$.
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Summary

- Smarr relation must be modified for $\Lambda \neq 0$.
- Thermodynamical volume: $V = \left. \frac{\partial M}{\partial P} \right|_{S,J,Q}$ ($P = -\frac{\Lambda}{8\pi}$).
- Reverse isoperimetric inequality, $V > V_0$, for AdS.
- Including pressure, first law becomes:

$$dU = TdS + \Omega dJ + \Phi dQ - PdV.$$

- Compressibility, $\kappa = -\frac{1}{V} \left. \frac{\partial V}{\partial P} \right|_{S,J} \geq 0$.
- Efficiency of Penrose process modified:
e.g. $Q = 0$: $\eta_{max} \approx 0.5184$ for an extremal black hole
($\eta \approx 0.2929$ without the PdV term).
- Generalises to higher dimensions.