

# Higgs phenomenology in the triplet extension of the MSSM

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Based on

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# Outline

- 1 Introduction
  - Motivation
  - The model
- 2 Higgs Phenomenology
  - Spectrum
  - Features of the decoupling limit
  - Features at small  $m_A$
- 3 Conclusion

# Motivation

- No clear discrepancies between data and SM predictions with  $m_h \simeq 126 \text{ GeV}$
- If we do not give up with the (Planck/GUT - EW) hierarchy problem, **SUSY** is (one of) the favourite UV option
- In the **MinimalSSM**  $m_h \simeq 126 \text{ GeV}$  requires “heavy” stop sector  $\Rightarrow$  **Little Hierarchy Problem**, i.e. some fine-tuning
- Non-minimalSSM models can alleviate this problem as they can enhance the tree-level Higgs mass via
  - D-terms: Extra gauge interactions
  - F-terms: Extra scalar sector (singlets and/or triplets)

If  $\Gamma(h \rightarrow \gamma\gamma)$  is confirmed by future data...

... the extra **charged** fermions in the triplet superfield are (potentially) largely coupled to the Higgs and can make the job

# The $Y=0$ Triplet Extension

$$\Sigma = \begin{pmatrix} \xi^0/\sqrt{2} & -\xi_2^+ \\ \xi_1^- & -\xi^0/\sqrt{2} \end{pmatrix}, \quad \Delta W = \lambda H_1 \cdot \Sigma H_2 + \frac{1}{2} \mu_\Sigma \text{tr} \Sigma^2 + \mu H_1 \cdot H_2$$

- $T$  parameter bound requires  $\langle \xi^0 \rangle \lesssim 4 \text{ GeV}$  which imposes (unless of fine-tuning)

$$|A_\lambda|, |\mu|, |\mu_\Sigma| \lesssim \frac{m_\Sigma^2 + \lambda^2 v^2 / 2}{10^2 \lambda v}$$

- This hierarchy implies decoupling between  $\xi^0$  and  $H_1, H_2$

$$\text{Mass boost: } V(H_1, H_2) \simeq V_{MSSM} + \lambda^2 |H_1^0 H_2^0|^2$$

(other tripl. SUSY extens. in E.J.Chun&al,1209.1303; Z.Kang&al,1301.2204)

# The relevant spectrum

- Heavy scalar triplet [  $\gtrsim 1 \text{ TeV}$  ]
- Minimiz. conditions

$$m_3^2 = m_A^2 \sin \beta \cos \beta, \quad m_Z^2 = \frac{m_2^2 - m_1^2}{\cos 2\beta} - m_A^2 + \frac{\lambda^2}{2} v^2$$

- CP-odd/charged Higgses [ no preferences ]

$$m_A^2 = m_1^2 + m_2^2 + 2|\mu|^2 + \frac{\lambda^2}{2} v^2, \quad m_{H^\pm}^2 = m_A^2 + m_W^2 + \frac{\lambda^2}{2} v^2$$

- Charginos [  $\mathcal{O}(100 \text{ GeV})$  if diphoton excess ]

$$\left( \tilde{W}^-, \tilde{H}_1^-, \tilde{\xi}_1^- \right) \mathcal{M}_{ch}^\pm \begin{pmatrix} \tilde{W}^+ \\ \tilde{H}_2^+ \\ \tilde{\xi}_2^+ \end{pmatrix}, \quad \mathcal{M}_{ch}^\pm = \begin{pmatrix} M_2 & gv_2 & 0 \\ gv_1 & \mu & -\lambda v_2 \\ 0 & -\lambda v_1 & \mu_\Sigma \end{pmatrix}$$

# The relevant spectrum

- CP-even Higgs masses (basis  $h_2, h_1$ )

$$\mathcal{M}_0^2 = \begin{pmatrix} m_A^2 \cos^2 \beta + m_Z^2 \sin^2 \beta & (\lambda^2 v^2 - m_A^2 - m_Z^2) \sin 2\beta/2 \\ (\lambda^2 v^2 - m_A^2 - m_Z^2) \sin 2\beta/2 & m_A^2 \sin^2 \beta + m_Z^2 \cos^2 \beta \end{pmatrix}$$

- After including radiative corrections  $\Delta\mathcal{M}_t^2(h_t)$  and  $\Delta\mathcal{M}_\Sigma^2(\lambda)$  (also in the min.condts) [ $m_h = 126 \text{ GeV}$ ]

$$\begin{pmatrix} h_2 \\ h_1 \end{pmatrix} = \begin{pmatrix} \cos \alpha & \sin \alpha \\ -\sin \alpha & \cos \alpha \end{pmatrix} \begin{pmatrix} h \\ H \end{pmatrix}$$

Coupling ratios  $r_{\mathcal{H}XX} = g_{\mathcal{H}XX}/g_{hXX}^{\text{SM}}$  ( $\mathcal{H} = h, H$ )

$r_{hVV}^0$	$r_{HVV}^0$	$r_{htt}^0$	$r_{Htt}^0$	$r_{hdd}^0$	$r_{Hdd}^0$
$\sin(\beta - \alpha)$	$\cos(\beta - \alpha)$	$\frac{\cos \alpha}{\sin \beta}$	$\frac{\sin \alpha}{\sin \beta}$	$-\frac{\sin \alpha}{\cos \beta}$	$\frac{\cos \alpha}{\cos \beta}$

# Decoupling limit $m_A \rightarrow \infty$

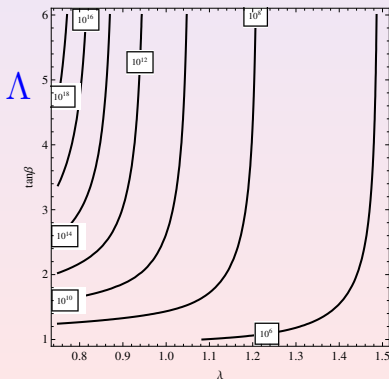
If  $m_A \gg m_h \dots$

Decoupling limit  $m_A \rightarrow \infty$ 

## Tree-level Higgs mass

$$m_h^2 = m_Z^2 \cos^2 2\beta + \lambda^2 v^2 \sin^2 2\beta / 2$$

Little hierarchy problem pushes towards small  $\tan \beta$  and large  $\lambda$ , but perturbation theory may break down at the scale  $\Lambda$



curves:  
 $10^{18}, 10^{16}, 10^{14}, \dots, 10^6$



Decoupling limit  $m_A \rightarrow \infty$ Signal strength  $\mathcal{R}_{\gamma\gamma} \quad [ = BR(h \rightarrow \gamma\gamma) / BR(h \rightarrow \gamma\gamma)_{SM} ]$ 

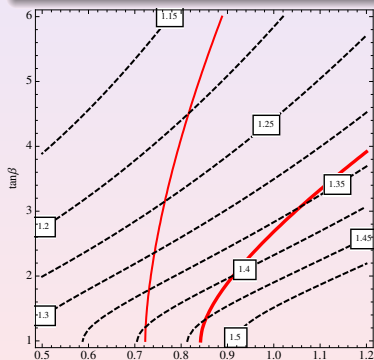
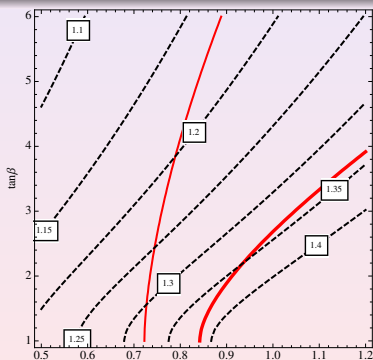
$$\frac{\partial}{\partial \log v} \log \det \mathcal{M}_{ch}(v) = - \frac{v^2 (\lambda^2 M_2 + g^2 \mu_\Sigma) \sin 2\beta}{M_2 \mu \mu_\Sigma - \frac{1}{2} \lambda^2 v^2 (\lambda^2 M_2 + g^2 \mu_\Sigma) \sin 2\beta}$$

- ATLAS:  $1.65 \pm 0.30$ ; CMS:  $0.8 - 1.1 \pm 0.3$
- Loop-induced process which is sensitive to new charged particles
- New triplet charged fermion can enhance  $R_{\gamma\gamma}$  ( $\lesssim 1.1$  via MSSM charginos; see Kraml's talk)
- As **no modification in the Higgs production** for  $m_A \rightarrow \infty$ , diphoton enhancement (easily implemented from the QED eff.pot. Ellis&al,79; Shifman&al,79; Carena&al,12)

$$R_{\gamma\gamma} = \left| 1 + \left( \frac{4}{3} \frac{\partial}{\partial \log v} \log \det \mathcal{M}_{ch}(v) \right) / \left( A_1(\tau_W) + \frac{4}{3} A_{1/2}(\tau_t) \right) \right|^2$$

Decoupling limit  $m_A \rightarrow \infty$ Signal strength  $\mathcal{R}_{\gamma\gamma}$   $[= BR(h \rightarrow \gamma\gamma)/BR(h \rightarrow \gamma\gamma)_{SM}]$ 

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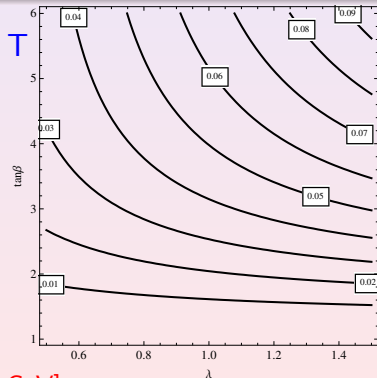
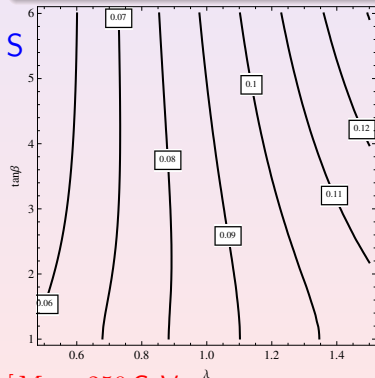
 $M_2 = 250 \text{ GeV}^\lambda$ 
 $X_t = 4, \mathbf{0}; m_Q = 700 \text{ GeV}$   
 $\mu = \mu_\Sigma \rightarrow m_{\chi_1^\pm} \sim 105 \text{ GeV}$ 
 $M_2 = 750 \text{ GeV}$

Decoupling limit  $m_A \rightarrow \infty$ 

S and T parameters (due to EWKinos)

$$\alpha S = \frac{s_W^2 \lambda^2}{10\pi^2} \frac{m_W^2}{\mu^2} \left[ 1 + \frac{19}{24} \sin 2\beta \right] + \mathcal{O}(g^4), \quad S = 0.04 \pm 0.09$$

$$\alpha T = \frac{3\lambda^2}{128\pi^2} \frac{m_W^2}{\mu^2} \cos^2 2\beta + \mathcal{O}(g^4), \quad T = 0.07 \pm 0.08$$



$[M_2 = 250 \text{ GeV}, \mu = \mu_\Sigma \rightarrow m_{\chi_1^\pm} \sim 105 \text{ GeV}]$

Small  $m_A$ 

And now let us

**DECREASE**  $m_A$

without upsetting the  $h$  Higgs  
signatures

Small  $m_A$ 

- CP-even Higgs masses (basis  $h_2, h_1$ )

$$\mathcal{M}_0^2 = \begin{pmatrix} m_A^2 \cos^2 \beta + m_Z^2 \sin^2 \beta & (\lambda^2 v^2 - m_A^2 - m_Z^2) \sin 2\beta/2 \\ (\lambda^2 v^2 - m_A^2 - m_Z^2) \sin 2\beta/2 & m_A^2 \sin^2 \beta + m_Z^2 \cos^2 \beta \end{pmatrix}$$

$$\begin{pmatrix} h_2 \\ h_1 \end{pmatrix} = \begin{pmatrix} \cos \alpha & \sin \alpha \\ -\sin \alpha & \cos \alpha \end{pmatrix} \begin{pmatrix} h \\ H \end{pmatrix}$$

- Tree-level  $h$  couplings are SM-like if  $\alpha = \beta - \pi/2$ . With  $\mathcal{M}_0^2$ :

$$\beta_c = \frac{\pi}{4}, \quad \lambda_c = \sqrt{2} \frac{m_h}{v}$$

- More generically with  $\mathcal{M}^2$ , fixing  $h$  mass yields (no  $m_A^4!$ ):

$$A(\tan \beta, \lambda, m_h) m_A^2 + B(\tan \beta, \lambda, m_h) = 0$$

SM-LIKE POINT  $(\lambda_c, \tan \beta_c)$  INDEPENDENT OF  $m_A$

Small  $m_A$ 

- CP-even Higgs masses (basis  $h_2, h_1$ )

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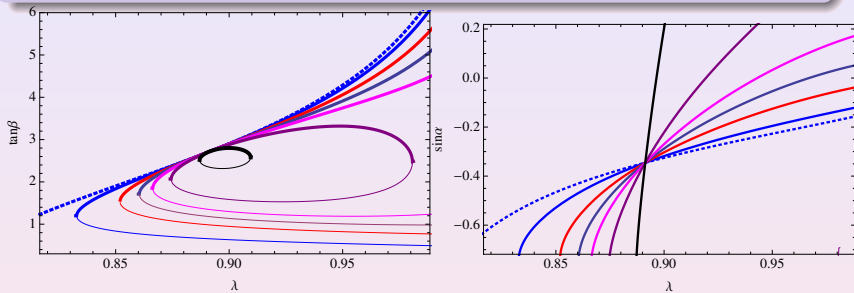
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SM-LIKE POINT  $(\lambda_c, \tan \beta_c)$  INDEPENDENT OF  $m_A$

Small  $m_A$ 

$\tan \beta(\lambda)$  and  $\sin \alpha(\lambda)$  fixing  $m_h = 126$  GeV



Radiative correction  $\Delta \mathcal{M}_t^2(h_t)$  and  $\Delta \mathcal{M}_\Sigma^2(\lambda)$  included  
 ( $m_Q = 700$  GeV,  $A_t = 0$ ,  $m_\Sigma = 5$  TeV).

Curves:  $m_A = \infty, 200, 155, 145, 140, 135, 130$  GeV

- No attempt to reproduce data with  $m_H = 126$  GeV (but worth to check!)

Small  $m_A$ Signal strengths  $\mathcal{R}_{\mathcal{H}XX}$ 

$$\mathcal{R}_{\mathcal{H}XX} = \frac{\sigma(pp \rightarrow \mathcal{H}) BR(\mathcal{H} \rightarrow XX)}{[\sigma(pp \rightarrow h) BR(h \rightarrow XX)]_{SM}}$$

$$\mathcal{R}_{\mathcal{H}XX}^{(ggF)} = \mathcal{R}_{\mathcal{H}XX}^{(\mathcal{H}tt)} = \frac{r_{\mathcal{H}tt}^2 r_{\mathcal{H}XX}^2}{\mathcal{D}}, \quad \mathcal{R}_{\mathcal{H}XX}^{(VBF)} = \mathcal{R}_{\mathcal{H}XX}^{(V\mathcal{H})} = \frac{r_{\mathcal{H}WW}^2 r_{\mathcal{H}XX}^2}{\mathcal{D}}$$

$$\begin{aligned} \mathcal{D} = & BR(h \rightarrow b b)_{SM} r_{hbb}^2 + BR(h \rightarrow gg, cc)_{SM} r_{\mathcal{H}tt}^2 \\ & + BR(h \rightarrow \tau\tau)_{SM} r_{\mathcal{H}\tau\tau}^2 + BR(h \rightarrow WW, ZZ)_{SM} r_{\mathcal{H}WW}^2 \end{aligned}$$

- No extra inv. width, i.e.  $m_{\chi_0} \gtrsim m_{\mathcal{H}}/2$
- sbottom-gluino may correct  $r_{hbb}$  ( $M_3 = 1 \text{ TeV}, m_{\tilde{b}} = 700 \text{ GeV}$ )

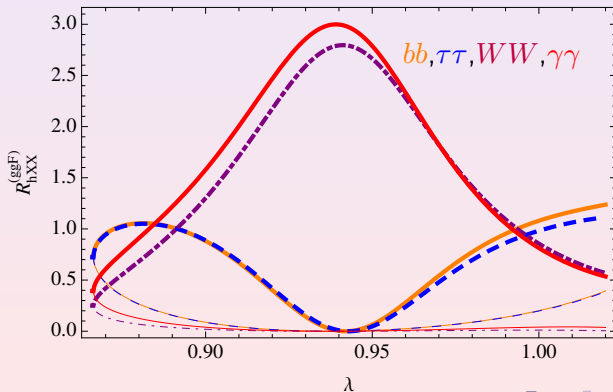


Small  $m_A$ 

$$m_A = 140 \text{ GeV}, \mu = \mu_\Sigma = 250 \text{ GeV}, m_{\chi^\pm} = 104 \text{ GeV}$$

Signal strengths  $\mathcal{R}_{hXX}$  from ggF and htt

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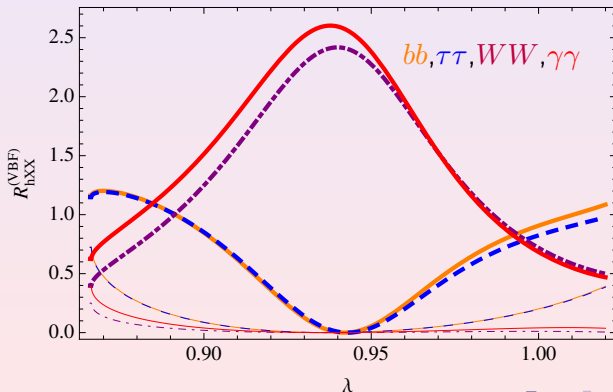


Small  $m_A$ 

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Signal strengths  $\mathcal{R}_{hXX}$  from VBF and Vh

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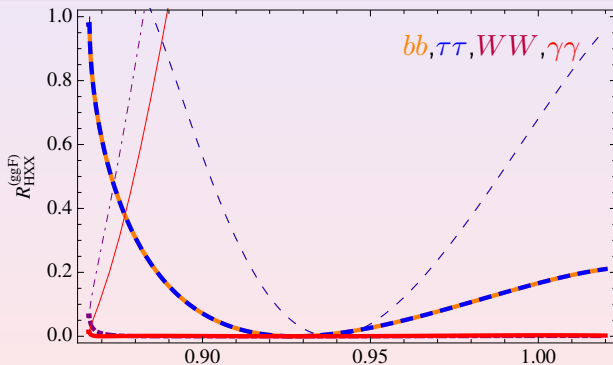


Small  $m_A$ 

$$m_A = 140 \text{ GeV}, \mu = \mu_\Sigma = 250 \text{ GeV}, m_{\chi^\pm} = 104 \text{ GeV}$$

Signal strengths  $\mathcal{R}_{HXX}$  from ggF and htt

$$\mathcal{R}_{\mathcal{H}XX} = \frac{\sigma(pp \rightarrow \mathcal{H})BR(\mathcal{H} \rightarrow XX)}{[\sigma(pp \rightarrow h)BR(h \rightarrow XX)]_{SM}}$$



- Further reduction if  $m_{\chi_0} \lesssim m_H/2$  ( $m_H \sim 138 \text{ GeV}$ )

# Conclusion

- Triplet extension alleviates the fine-tuning with respect to the MSSM
- At very small  $m_A$  there are 2 regions around  $\lambda \approx 0.9$  and  $\lambda \approx 1$  that mimic the signal strength of a SM Higgs.
- Some small deviations from a pure SM Higgs (e.g.  $\gamma\gamma$ ,  $bb$ ,  $\tau\tau$ ) can also be encompassed without need of modifying other rates
- The extra Higgs sector may be at the EW scale but hidden
- More data and analyses on  $H^\pm$  and  $A$  are worthwhile
- Good alternative to singlet extensions

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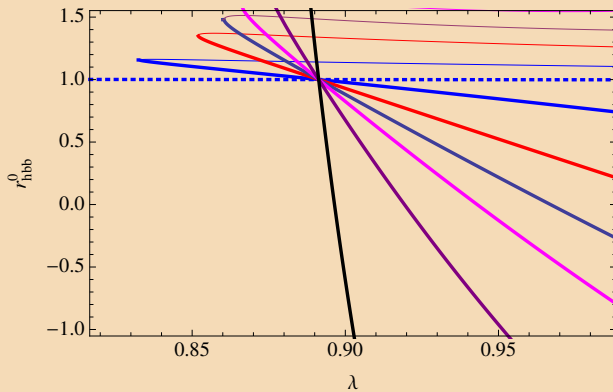
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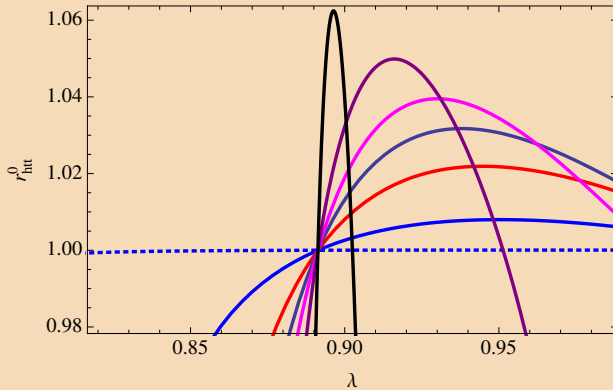


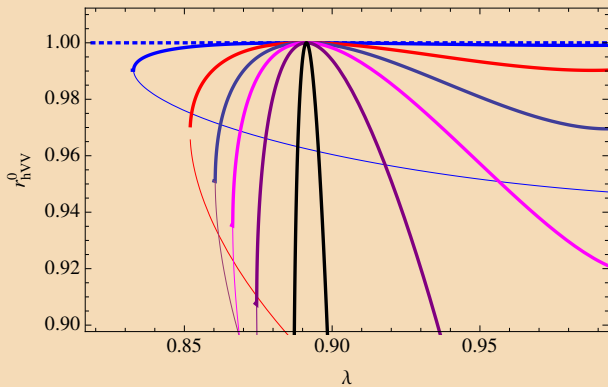
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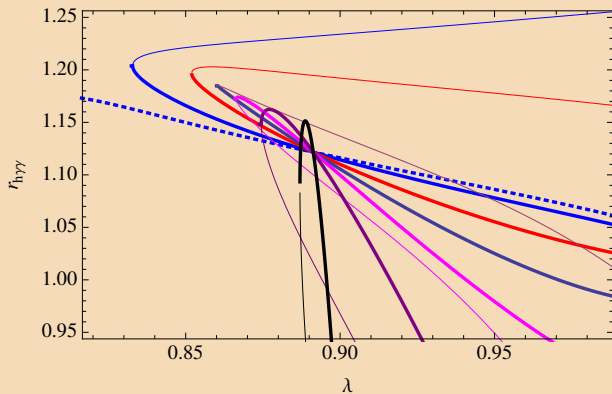
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# EXTRAS

Plot of  $r_{hdd}^0$ 

Plot of  $r_{\text{Htt}}^0$ 

Plot of  $r_{HVV}^0$ 

Plot of  $r_{h\gamma\gamma}$ 

$m_{\chi_1^\pm} = 104$  GeV (solid line),  $m_{\chi_1^\pm} = 150$  GeV with  $\mu = \mu_\Sigma = 300$  GeV (dashed line) and  $m_{\chi_1^\pm} = 200$  GeV with  $\mu = \mu_\Sigma = 350$  GeV (dotted line)

