

Bulk Scalar Phenomenology in Randall-Sundrum

arXiv: 1306.xxxx

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Outline

Introduction

Motivation

Scalars in Randall-Sundrum

Higgs-Radion Mixing

Mixing Term

Higgs and Radion Couplings

Results

General

Higgs

Radions

Conclusions

The Precision Higgs Era



Higgs boson-like particle discovery claimed at LHC



TIME Talks to the Physicists Who Found the Higgs

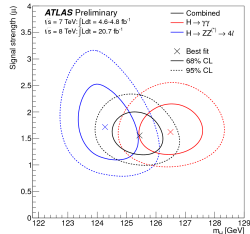
By Jeffrey Kluger Thursday, July 19, 2012



- LHC discovery last July
 - Did you notice?
- Many Higgs couplings observed
- New era of particle physics

In particular:

- Higgs mass at percent level
- Coupling uncertainties shrinking
- Is everything post-GSW irrelevant?



Higgs Constraining BSM

Higgs measurements constraining deviations from SM:

Carmi *et al* 1202.3144, Low *et al* 1207.1093, Giardino *et al* 1207.1347, Ellis & You 1207.1693, Bonnet *et al* 1207.4599, Plehn & Rauch 1207.6108, Cacciapaglia *et al* 1210.8120, Belanger *et al* 1302.5694 **And many more...**

What distinguishes us?

- Specific, top-down model: Randall-Sundrum
- Kinetic **and** mass mixing from single coupling
- Interesting effects on Radion phenomenology
- Relate to full RS parameter space

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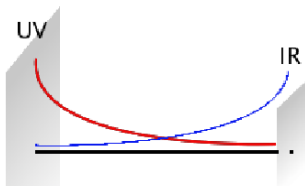
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Randall-Sundrum Models

Brief Review



Warped Extra Dimensions:

- Popular BSM framework
- Model of Flavour Physics
- Solution to hierarchy problem:

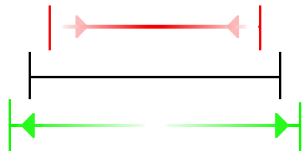
- Non-factorisable warping in fifth dimension
- Gravitational redshift exponentially lowers scales

$$ds^2 = e^{-2ky} \eta_{\mu\nu} dx^\mu dx^\nu - dy^2 \quad \Rightarrow \quad \Lambda_{IR} \sim \Lambda_{UV} e^{-kL}$$

AdS-CFT: IR-localised Higgs \sim 4D composite Higgs

The Radion

- **New scalar mode** in all RS models
- Fluctuations in **size** of extra dim.
(See cartoon \Rightarrow)

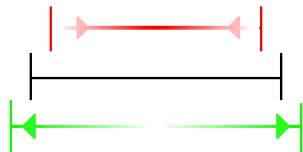


$$ds^2 = e^{-2ky-2G(x,y)} dx^\mu dx_\mu - (1 + 2G(x,y))^2 dy^2$$

- Typically very light, $m_r \ll m_{KK}$
 - Massless without stabilisation! *e.g.* Goldberger & Wise hep-ph/9907447
 - Neglect back-reaction in this talk: $G(x,y) \sim e^{2ky}$
 - Can be only new state observed
- Higgs-Radion mixing will be generic in RS models

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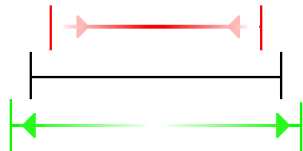


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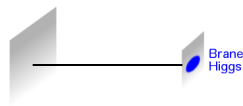
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The Brane Higgs

Review

Usual RS approach: Higgs on brane

$$\int d^5x \sqrt{|g|} \delta(y - L) [(D^\mu H)^\dagger D_\mu H - V(H)]$$



Canonical Higgs rescaling: Solution to Hierarachy Problem

$$H_{phys} = e^{kL} H \quad \Rightarrow \quad m_{4D}^2 = e^{-2kL} m_{5D}^2$$

SM masses brane-localised: Higgs couplings as SM

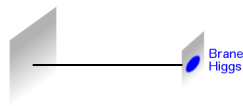
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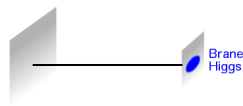
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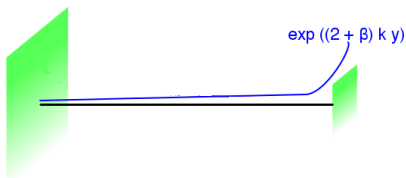
The Bulk Higgs

Alternative: Higgs in 5D Bulk.

$$\int d^5x \sqrt{|g|} [(D^M H)^\dagger D_M H - V_{bk}(H) - V_{br}(H)]$$



Break symmetry **only** with IR potential: Solve hierarchy problem



- $\beta^2 = 4 + m_H^2/k^2$
- $v(y) \propto e^{(2+\beta)ky}$
- $h^0(y) \propto v(y) + \mathcal{O}\left(\frac{m_h^2}{m_{KK}^2}\right)$

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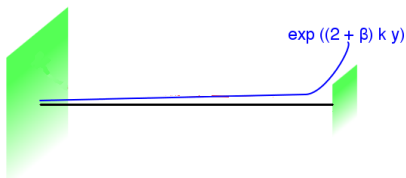
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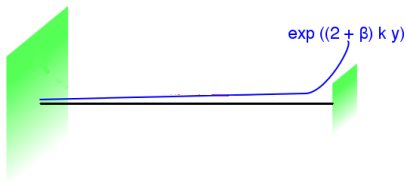
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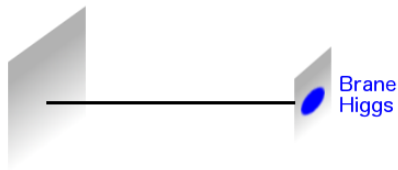


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SM masses in bulk: **Higgs couplings as SM**

Bulk vs Brane Higgs

Brane Higgs:



- More predictive
- More studied
- Couplings as SM
- **Kinetic** Mixing with Radion
- Usual Radion couplings

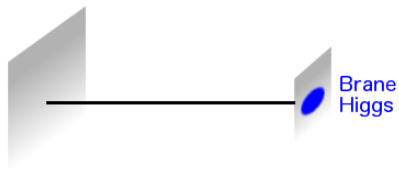
Bulk Higgs:



- Generalisation of model
- More room to explore
- Couplings as SM
- **Kinetic and Mass** Mixing
- No Radion- WW coupling

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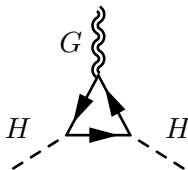
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Origin of the Mixing

Radion \subset metric: unique renormalisable mixing term

$$\mathcal{L}_{mix} = \xi R H^\dagger H$$

- Studied in 4D;
 $\xi = \frac{1}{6}$ **required** for $T_\mu^\mu = 0$
- Generated by top (+ ...) loops
- RS context:
 Exists both in bulk and on brane



$$\sim -i \frac{y_t^2}{4\pi} \frac{m_t}{\Lambda_G}$$

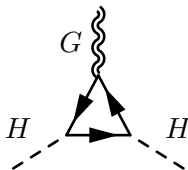
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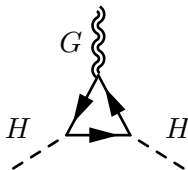
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$$\sim -i \frac{y_t^2}{4\pi} \frac{m_t}{\Lambda_G}$$

Treat **size and sign** of coupling as **free parameter**

Brane Mixing

Kinetic Mixing

On brane, must use induced metric, curvature

$$\begin{aligned} \mathfrak{g}_{\mu\nu}(x) &= g_{\mu\nu}(x, L) & R_4 &\equiv R(\mathfrak{g}) \\ &= e^{-2kL} \eta_{\mu\nu} (1 - 2G) + \dots & &= -6 e^{2kL} \square G + \dots \end{aligned}$$

- Induced mixing is **pure kinetic**

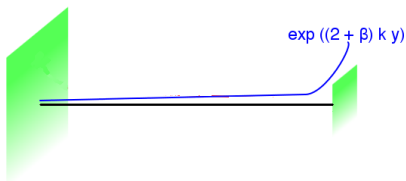
$$\xi R_4 H^\dagger H \supset \sqrt{\frac{2}{3}} \frac{v}{M_{Pl} e^{-kL}} \xi (\partial^\mu r)(\partial_\mu h^0)$$

Bulk Mixing I

Fiddling with the Profile

- $R \neq 0$ so mixing term \Rightarrow Higgs Mass

$$-c_H^2 k^2 H^\dagger H + \xi R_5 H^\dagger H = -(c_H^2 + 20\xi) k^2 H^\dagger H + \dots$$



- Shift: $\beta^2 = 4 + c_H^2 + 20\xi$
- VEV, Higgs profile **unchanged** as functions of β
- Note: mass parameter varies

Bulk Mixing II

Kinetic and Mass Mixing

Bulk mixing linear terms **fun!**

$$R_5 = \dots + R^{MN}(\delta g)_{MN} + \nabla^M \nabla^N (\delta g)_{MN} - \nabla_M \nabla^M (\delta g)_N{}^N + \dots$$

Kinetic Mixing:

$$\sqrt{\frac{2}{3}} \frac{1 + \beta}{2 + \beta} \frac{v}{M_{Pl} e^{-kL}} \xi (\partial^\mu r)(\partial_\mu h^0)$$

Mass Mixing: (several contributions)

$$\sqrt{\frac{2}{3}} 2(1+4\beta) \frac{v}{M_{Pl} e^{-kL}} (k e^{-kL})^2 \xi r(x) h^0(x)$$

- Parametric suppression
- β -dependence
- Large mass scale
 $k e^{-kL} \gg m_r, m_h$

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Diagonalisation

See Any Field Theory Textbook

Effective 4D theory:

$$\mathcal{L}_{mix} = \frac{1}{2}(\partial h^0)^2 + \xi\alpha(\partial h^0)(\partial r) + \frac{1}{2}(\partial r)^2 - \frac{1}{2}m_h^2(h^0)^2 - \xi\gamma\Lambda_{IR}^2 h r - \frac{1}{2}m_r^2 r^2$$

$$\alpha, \gamma \sim \frac{v}{M_{Pl} e^{-kL}}$$

$$\Lambda_{IR} = k e^{-kL}$$

$$\gamma = 0 \text{ (brane)}$$

- Kinetic mixing: **Non-unitary transformation**

$$r \rightarrow \tilde{r} - \xi \alpha h^0 + \mathcal{O}(\alpha^2)$$

- Mass mixing: **Orthogonal rotation**

$$\begin{pmatrix} r_m \\ h_m^0 \end{pmatrix} = \begin{pmatrix} c_\theta & s_\theta \\ -s_\theta & c_\theta \end{pmatrix} \begin{pmatrix} \tilde{r} \\ h^0 \end{pmatrix}, \quad s_\theta = 2\xi\gamma\Lambda_{IR}^2/m_r^2 + \mathcal{O}(\gamma^2)$$

- Note: linear approximation here for illustration **only**

Modified Couplings

In linear approximation

$$\begin{aligned}
 -\mathcal{L}_{int} &= h^0 \mathcal{O}_m + r \mathcal{O}_r \\
 &\rightarrow h_m^0 \left[\mathcal{O}_m - \xi \left(\alpha + \gamma \frac{\Lambda_{IR}^2}{m_r^2} \right) \mathcal{O}_r \right] + r_m \left[\mathcal{O}_r + \xi \gamma \frac{\Lambda_{IR}^2}{m_r^2} \mathcal{O}_h \right]
 \end{aligned}$$

- **Higgs couplings** contribution from both mixings
 - Couplings enhanced **or** suppressed (Sign of ξ)
- Radion couplings contribution from mass mixing only

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Radion Couplings

Relevant interactions: r -SM-SM

1. Kinetic terms:

- **Same** for bulk and brane Higgs
- Contribute to **gluon and photon** couplings:

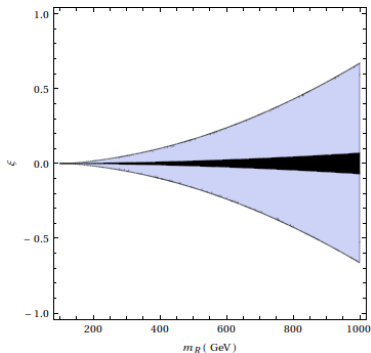
$$\mathcal{L}_{int} \supset \frac{1}{8} F^{\mu\nu} F_{\mu\nu} r(x) \frac{1}{kL} \frac{1}{\sqrt{3}M_{Pl}e^{-kL}}$$

2. Mass terms:

- **Different** for bulk and brane Higgs
- Bulk couplings β and $c_{L,R}$ dependent
- Coupling to W, Z **vanishes** for bulk Higgs

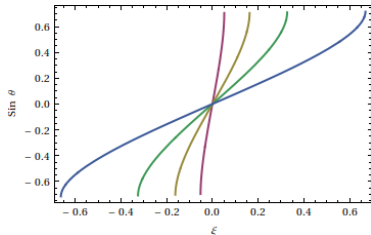
$$\begin{aligned} \sqrt{|g|} g^{\mu\nu} &= \left[e^{-4kL} (1 - 2G)^2 (1 + 2G) \right] \left[e^{2kL} \eta^{\mu\nu} (1 + 2G) \right] \\ &= e^{-2kL} \eta^{\mu\nu} + \mathcal{O}(G^2) \end{aligned}$$

Theoretical Mixing Limit



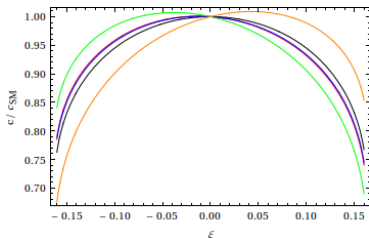
- Mixing increases radion mass
- Radion mass as input \Rightarrow upper limit on $|\xi|$
- Shown: $\beta = 1$, $\Lambda_{IR} = 1$ TeV
- ($\Lambda_{IR} = 10$ TeV in black)
- Brane Higgs: **much** weaker limit

Scalar Mixing Angle



- Despite small ξ , large angle θ
- Shown: $\beta = 1, \Lambda_{IR} = 1 \text{ TeV}$
- $m_r = 200, 300, 500, 1000 \text{ GeV}$ for red, yellow, green, blue lines

Higgs Couplings



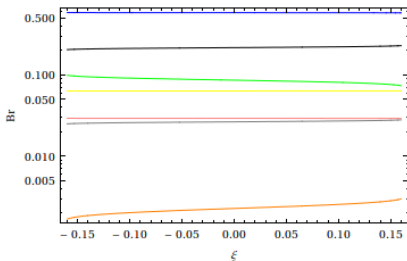
- Note asymmetry in ξ
- Mass mixing dominant; suppressing all couplings
- Kinetic mixing relevant for small $|\xi|$
- γ enhanced when other couplings suppressed

On same scale, brane Higgs couplings unchanged

($\beta = 1, \Lambda_{IR} = 1 \text{ TeV}, m_r = 500 \text{ GeV}$)

Black: W/Z ; Red/Blue: t, b ; Green: g ; Orange: γ

Higgs Partial Widths



- Largest effects on g, γ
- Asymmetry in ξ

$(\beta = 1, \Lambda_{IR} = 1 \text{ TeV}, m_r = 500 \text{ GeV})$

Black: W/Z ; Blue: b ; Yellow: τ ; Green: g ; Red: c ; Orange: γ

Higgs Searches

- Use most recent (post-*Moriond*) public results
- Observations: signal normalised to SM expectation, μ :

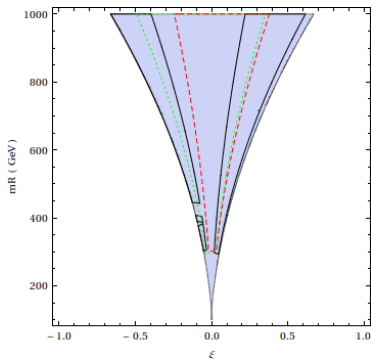
$$\mu_X = (a|c_g|^2 + b|c_W|^2) \frac{\Gamma(h \rightarrow X)}{\Gamma^{SM}(h \rightarrow X)}$$

for process-dependent a, b

- Combine ATLAS, CMS in simple fashion: weighted mean

	WW	ZZ	$\gamma\gamma$	bb	$\tau^+\tau^-$
ATLAS	1.0 ± 0.3	1.5 ± 0.4	1.6 ± 0.3	-0.4 ± 1.0	0.8 ± 0.7
CMS	0.68 ± 0.2	0.92 ± 0.21	0.77 ± 0.27	1.15 ± 0.62	1.1 ± 0.4
Average	0.78 ± 0.17	1.0 ± 0.2	1.1 ± 0.2	0.72 ± 0.53	1.0 ± 0.4

Higgs Constraints



- Shaded: theory allowed
- No constraints from $h \rightarrow f\bar{f}$
- Contours: 1- σ allowed regions
- On same scale, brane Higgs couplings identical to SM

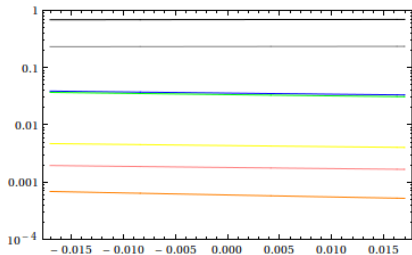
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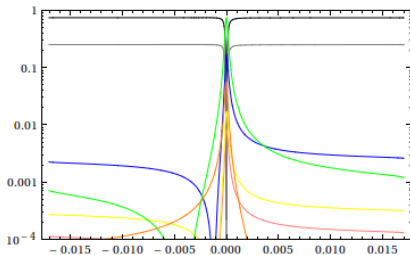
Radion Partial Widths

200 GeV

Brane Higgs:



Bulk Higgs:



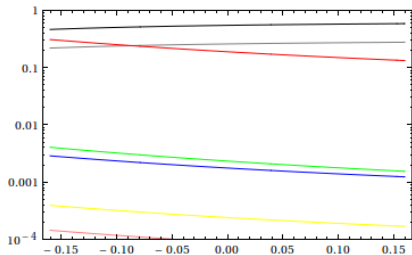
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Bulk Higgs suppresses width to WW, ZZ

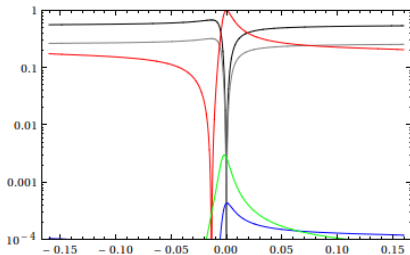
Radion Partial Widths

500 GeV

Brane Higgs:



Bulk Higgs:



Black: W/Z ; Red: t ; Blue: b ; Yellow: τ ; Green: g

Interesting $t\bar{t}$ signals?

Conclusions & Future Directions

1. Radion and bulk Higgs have both **mass** and **kinetic** mixing
2. **Mass mixing dominates** phenomenology, set by KK scale
3. Have used Higgs measurements to **constrain mixing ξ**
4. Have shown how **radion couplings differ** for bulk Higgs
5. To do: use LHC results to **constrain radion** parameter space